

1           **Resolving weather fronts increases the large-scale**  
2           **circulation response to Gulf Stream SST anomalies in**  
3           **variable-resolution CESM2 simulations**

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9           **Key Points:**

- 10           • There is a large circulation response to idealized Gulf Stream SST anomalies in  
11           an atmospheric model with 14-km regional grid refinement
- 12           • This response is weaker or absent in simulations with 28-km or coarser resolution,  
13           which do not fully resolve mesoscale frontal processes
- 14           • Transient-eddy fluxes of heat and momentum are modified as fronts pass over warm  
15           SSTs, leading to a large-scale circulation response

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**Abstract**

Canonical understanding based on general circulation models (GCMs) is that the atmospheric circulation response to midlatitude sea-surface temperature (SST) anomalies is weak compared to the larger influence of tropical SST anomalies. However, the  $\sim 100$ -km horizontal resolution of modern GCMs is too coarse to resolve strong updrafts within weather fronts, which could provide a pathway for surface anomalies to be communicated aloft. Here, we investigate the large-scale atmospheric circulation response to idealized Gulf Stream SST anomalies in Community Atmosphere Model (CAM6) simulations with 14-km regional grid refinement over the North Atlantic, and compare it to the responses in simulations with 28-km regional refinement and uniform 111-km resolution. The highest resolution simulations show a large positive response of the wintertime North Atlantic Oscillation (NAO) to positive SST anomalies in the Gulf Stream, a 0.4-standard-deviation anomaly in the seasonal-mean NAO for  $2^\circ\text{C}$  SST anomalies. The lower-resolution simulations show a weaker response with a different spatial structure. The enhanced large-scale circulation response results from an increase in resolved vertical motions with resolution and an associated increase in the influence of SST anomalies on transient-eddy heat and momentum fluxes in the free troposphere. In response to positive SST anomalies, these processes lead to a stronger and less variable North Atlantic jet, as is characteristic of positive NAO anomalies. Our results suggest that the atmosphere responds differently to midlatitude SST anomalies in higher-resolution models and that regional refinement in key regions offers a potential pathway to improve multi-year regional climate predictions based on midlatitude SSTs.

**Plain Language Summary**

Variations in the ocean surface temperature (SST) influence the atmospheric circulation and thus climate over land. Canonical understanding is that tropical SSTs are more important than SSTs in midlatitudes. However, this understanding is based on climate models that don't resolve processes at scales less than 100 km. Here, we show that by increasing the atmospheric model resolution to resolve features on smaller scales, such as weather fronts, we find a larger atmospheric circulation response to midlatitude SST anomalies in the North Atlantic. North Atlantic SST anomalies can be predicted multiple years in advance, and a larger atmospheric circulation response to these predictable SST anomalies therefore implies increased predictability of climate over the surrounding land regions.

**1 Introduction**

North Atlantic sea-surface temperatures (SSTs) exhibit variability on seasonal to decadal timescales (e.g., Deser & Blackmon, 1993; R. Zhang et al., 2019), providing a potential source of predictability for atmospheric circulation and regional climate on these timescales. Recent work has improved our understanding of the ocean-atmosphere mechanisms governing North Atlantic SST variability (Menary et al., 2015; Delworth et al., 2017; S. Yeager & Robson, 2017; R. C. J. Wills et al., 2019; R. Zhang et al., 2019; Årthun et al., 2021) and shown that initialized climate models have skill in predicting the decadal evolution of North Atlantic SST (Msadek et al., 2014; Meehl et al., 2014; S. G. Yeager et al., 2018; D. Smith et al., 2019; Borchert et al., 2021; S. G. Yeager et al., 2023), but this will only help to make model-based predictions of regional climate anomalies in the surrounding continents if the models correctly simulate the atmospheric response to North Atlantic SST anomalies.

There is a large literature that tries to diagnose the atmospheric circulation response to North Atlantic SST anomalies from observations (see, e.g., Czaja & Frankignoul, 1999; Frankignoul et al., 2001; Czaja & Frankignoul, 2002; Gastineau et al., 2013; Gastineau & Frankignoul, 2015; S. M. Wills et al., 2016). However, the North Atlantic atmospheric

66 circulation exhibits strong internal variability, particularly due to the North Atlantic Os-  
 67 cillation (NAO), and this internal variability leads to a large intrinsic uncertainty in the  
 68 diagnosed relationship between SSTs and circulation. The relationship between SSTs  
 69 and circulation can be accurately diagnosed in climate model ensembles by averaging the  
 70 relationship over a large number of simulations with different initial conditions, but the  
 71 modeled relationship may not accurately reflect the real-world relationship. Indeed, while  
 72 the canonical understanding based on climate models is that the large-scale circulation  
 73 responds only weakly to midlatitude SST anomalies (Lau & Nath, 1994; Kushnir et al.,  
 74 2002), there is growing evidence that the atmospheric response to midlatitude SST anoma-  
 75 lies is systematically underestimated in climate models (Simpson et al., 2018, 2019; R. C. J. Wills  
 76 et al., 2019; Czaja et al., 2019), and that this may be rectified by increasing the atmo-  
 77 spheric resolution to resolve mesoscale processes over ocean frontal zones (Smirnov et  
 78 al., 2015; Sheldon et al., 2017; Czaja et al., 2019; Oldenburg et al., 2022; Famooss Paolini  
 79 et al., 2022; Seo et al., 2023).

80 Global climate models (GCMs) are typically run with  $\sim 100$  km or coarser horizon-  
 81 tal resolution and are therefore unable to simulate mesoscale atmospheric processes such  
 82 as the conditional symmetric instability and other frontal dynamics ( $\sim 10$ - $100$  km scales),  
 83 which are important in the dynamics of weather. Increasing atmospheric model resolu-  
 84 tion is known to increase the strength of resolved updrafts (Jeevanjee & Romps, 2016;  
 85 Herrington & Reed, 2018, 2020), including the ascent within weather fronts passing over  
 86 the Gulf Stream SST front (Sheldon et al., 2017). It is well established that a deeper mid-  
 87 latitude heating anomaly will induce a larger downstream circulation response (Hoskins  
 88 & Karoly, 1981), so if weather fronts could facilitate upward heat transport, this would  
 89 lead to a larger circulation response. However, the influence of resolving updrafts on large-  
 90 scale atmosphere-ocean coupling and seasonal-to-decadal predictability has not been sys-  
 91 tematically studied in climate models. A key factor limiting understanding is that cur-  
 92 rent global high-resolution atmospheric modeling efforts on climate timescales (i.e., run  
 93 for at least 10 years) are generally limited to  $1/4^\circ$  ( $\sim 25$  km) atmospheric resolution (Bacmeister  
 94 et al., 2014; Haarsma et al., 2016; Chang et al., 2020), which is still too coarse to fully  
 95 resolve weather fronts. It is extremely costly to run global models at sub-25-km atmo-  
 96 spheric resolution for the multiple decades needed to evaluate potential increases in the  
 97 circulation response to midlatitude SST anomalies and in predictability at seasonal-to-  
 98 decadal timescales.

99 In this work, we use variable-resolution (VR) simulations, where horizontal reso-  
 100 lution is enhanced only in a region of interest, to evaluate the potential benefit of resolv-  
 101 ing mesoscale processes for atmospheric predictability stemming from persistent mid-  
 102 latitude SSTs. VR modeling is widely used in weather forecasting (e.g., Buizza et al. (2007)),  
 103 but it is only starting to be explored for simulating climate variability and change (e.g.,  
 104 Zarzycki & Jablonowski, 2014; Zarzycki et al., 2015; van Kampenhout et al., 2019; Her-  
 105 rington et al., 2022; Wijngaard et al., 2023; Schemm, 2023). Here, we use VR configu-  
 106 rations of the spectral element dynamical core in the Community Atmosphere Model (CAM-  
 107 SE; P. H. Lauritzen et al., 2018), with 14-km ( $\sim 1/8^\circ$ ) or 28-km ( $\sim 1/4^\circ$ ) resolution over  
 108 the North Atlantic and Europe (Fig. 1), to model the large-scale atmospheric circula-  
 109 tion response to idealized North Atlantic SST anomalies. More details of the model and  
 110 grid configuration are provided in Sections 2.1 and 2.2, respectively.

111 In this paper, we focus on simulations with idealized SST anomalies in the Gulf  
 112 Stream region (Fig. 2; more details in Section 2.3). The Gulf Stream region is chosen  
 113 due to the large magnitude of observed SST variability in this region (S. M. Wills et al.,  
 114 2016) and the number of previous idealized modeling studies focusing on this region (Kaspi  
 115 & Schneider, 2011; Kuwano-Yoshida et al., 2010; O'Reilly et al., 2017; Sheldon et al., 2017).  
 116 Importantly, we use the same  $1^\circ$  resolution SSTs in all simulations, such that differences  
 117 in the atmospheric response between grids are only due to differences in atmospheric res-  
 118 olution. There is an extensive literature documenting how climatological SST biases (Chang

119 et al., 2020; Athanasiadis et al., 2022; Oldenburg et al., 2022) and boundary layer pro-  
 120 cesses over midlatitude fronts (Small et al., 2014; Seo et al., 2023) improve with ocean-  
 121 model resolution. We leave aside the important influence of ocean resolution for this study  
 122 in order to isolate the influence of atmospheric resolution. Follow up work should inves-  
 123 tigate how simultaneously resolving mesoscale processes in the atmosphere and ocean  
 124 influences the simulation of large-scale atmosphere-ocean coupling.

125 The paper is organized as follows. Details of the model used, the new variable-resolution  
 126 grids, and the idealized SST anomaly simulations are described in Section 2. The results  
 127 of these simulations are shown in Section 3, including subsections on the large-scale cir-  
 128 culation response, the projection of the response onto modes of internal variability, the  
 129 local air-sea interactions and cross-front circulation response, a thermodynamic analy-  
 130 sis, and the modification of transient eddy fluxes by the SST forcing. In Section 4, we  
 131 summarize our findings and discuss the implications for the signal-to-noise paradox and  
 132 seasonal-to-decadal predictability.

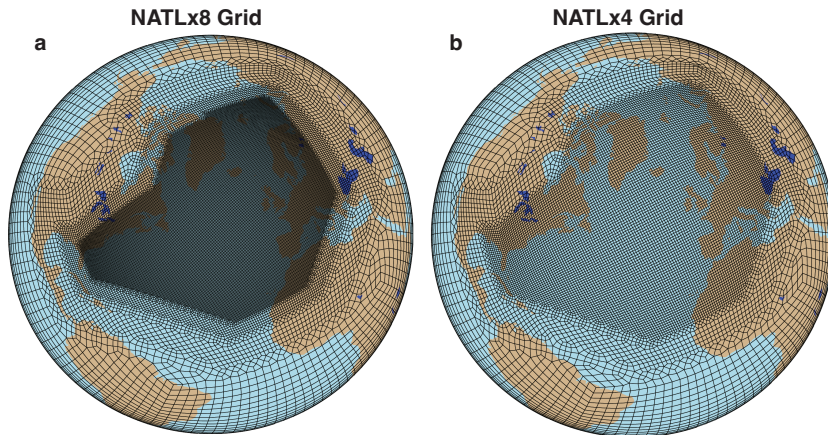
## 133 2 Variable-Resolution Simulations

### 134 2.1 Modeling Setup

135 Our simulations use the Community Earth System Model version 2.2 (Danabasoglu  
 136 et al., 2020; Herrington et al., 2022). Specifically, they use the Community Atmospheric  
 137 Model version 6 (CAM6), with the spectral element (SE) dynamical core (P. H. Lauritzen  
 138 et al., 2018), coupled to a data ocean (specified SST and sea ice) and the Community  
 139 Land Model version 5 (Lawrence et al., 2019). The atmosphere has 32 hybrid pressure-  
 140 sigma levels in all simulations, with a model top at  $\sim 2$  hPa.

141 The CAM6 physical parameterization package (Gettelman et al., 2019) contains  
 142 a high-order turbulence closure, Cloud Layers Unified By Binormals (CLUBB; Golaz et  
 143 al., 2002; Bogenschutz et al., 2013), which serves as a boundary layer, shallow convec-  
 144 tion and cloud macrophysics scheme. CLUBB is sub-cycled with a two-moment cloud  
 145 microphysics scheme (Gettelman & Morrison, 2015; Gettelman et al., 2015) and aerosol  
 146 activation scheme (Liu et al., 2007) for simulating cloud-aerosol interactions and precip-  
 147 itation processes. Deep convection is parameterized using a convective quasi-equilibrium  
 148 mass flux scheme (G. Zhang & McFarlane, 1995; Neale et al., 2008), supporting down-  
 149 drafts and convective momentum transport (Richter & Rasch, 2008). Boundary layer  
 150 form drag is parameterized after Beljaars et al. (2004) and orographic gravity waves are  
 151 parameterized using an anisotropic scheme that utilizes sub-grid orientations of ridges  
 152 derived from a high-resolution gridded topography data set (Danielson & Gesch, 2011).

153 The SE dynamical core is based on a cube-sphere grid, tiled with quadrilateral finite-  
 154 elements. The hydrostatic primitive equations are solved using the continuous-Galerkin  
 155 method (Taylor et al., 1997; Taylor & Fournier, 2010), with each element containing a  
 156 2D fourth-order polynomial basis set, and with  $4 \times 4$  quadrature nodes (i.e., grid points)  
 157 located at the roots of the basis functions. Grid points located on the element bound-  
 158 aries are shared with adjacent elements, facilitating communication between elements  
 159 via the direct stiffness summation (Canuto et al., 2007), and resulting in  $3 \times 3$  indepen-  
 160 dent grid points per element. For quasi-uniform grids, the SE method for tracer trans-  
 161 port is replaced with the Conservative Semi-Lagrangian Multi-tracer transport scheme  
 162 (CSLAM; P. H. Lauritzen et al., 2017), which operates on a separate finite-volume grid  
 163 containing  $3 \times 3$  control volumes per element. The physical parameterizations (hereafter  
 164 physics) are evaluated on the finite-volume grid in CSLAM, whereas in standard SE the  
 165 physics are evaluated at the quadrature points. A vertically Lagrangian scheme is used  
 166 in the vertical (Lin, 2004), wherein the 2D dynamics evolve in floating Lagrangian lay-  
 167 ers and are subsequently mapped back to a fixed Eulerian vertical grid.



**Figure 1.** Variable-resolution North Atlantic grids for CAM-SE: (a) The NATLx8 grid, with horizontal resolution varying from 14 km resolution in the North Atlantic to 111 km in the far field; (b) the NATLx4 grid, with horizontal resolution varying from 28 km resolution in the North Atlantic to 111 km in the far field. Note that what is shown is the element grid; the computational grid has  $3 \times 3$  independent grid points per element.

168 The SE dynamical core also supports variable-resolution grids, through invoking  
 169 scale-aware hyper-viscosity (Guba et al., 2014) and imposing rougher terrain in the re-  
 170 fined region, generated using CESM’s topography generation software (P. Lauritzen et  
 171 al., 2015). Variable-resolution currently does not support CSLAM, and the SE method  
 172 is used for tracer transport instead. The parameterizations are otherwise unmodified as  
 173 the refinement is increased. Notably, the deep convective parameterization is still included  
 174 for the maximum refinement used in this study (14-km grid spacing in refinement region),  
 175 though the convection scheme is known to become less active when the resolution is in-  
 176 creased and the physics time-step is reduced (Williamson, 2013; Herrington & Reed, 2020).  
 177 The SE time-stepping is reduced to satisfy the Courant-Friedrich-Lewy (CFL) condition  
 178 in the refined region, whereas the time-stepping in the physics is reduced to avoid large  
 179 time-truncation errors (Herrington & Reed, 2018). The physics time steps used are tab-  
 180 ulated based on the grid spacing of the refinement region in Herrington et al. (2022).

## 181 2.2 North Atlantic Variable-Resolution Grids and Performance

182 The basis for our regionally refined grids is the quasi-uniform ne30pg3 grid (here-  
 183 after NE30), which has  $30 \times 30$  quadrilateral elements per cubed sphere face and  $3 \times 3$   
 184 control volumes per element, for a total of 48,600 control volumes and an average hor-  
 185 izontal grid spacing of 111 km.

186 The North Atlantic (NATL) grids were generated using the software package SQuad-  
 187 gen (<https://github.com/ClimateGlobalChange/squadgen>) by rotating the cubed sphere  
 188 to have a face in the center of the North Atlantic, then refining a region mostly within  
 189 that face but extending also to neighboring faces (due to the irregular shape of the North  
 190 Atlantic). The NATLx8 grid has a maximum of  $8 \times$  refinement, i.e.,  $8 \times 8$  elements in  
 191 place of a single element in NE30, corresponding to a horizontal grid spacing of 14 km.  
 192 This refinement takes places in 3 steps, with  $2 \times$  and  $4 \times$  refinement regions for transi-  
 193 tion between the  $1 \times$  region and the  $8 \times$  region. The NATLx4 grid simply replaces all  $8 \times$   
 194 regions with  $4 \times$  refinement, corresponding to a horizontal grid spacing of 28 km. The  
 195 NATLx8 and NATLx4 grids have 317,567 and 142,346 control volumes, respectively.

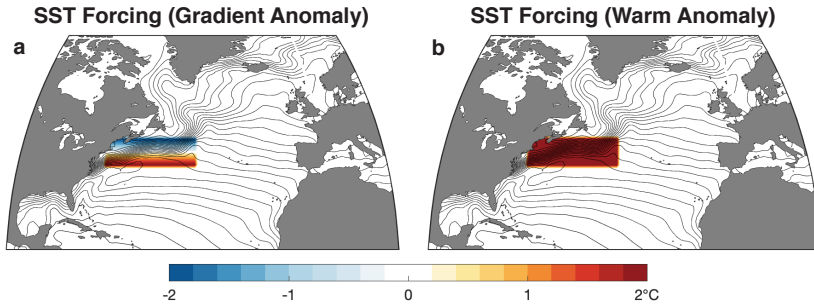
196 The refinement region for our simulations includes the Gulf Stream, which is the  
 197 primary region of focus for this work, but also extends to other regions of the North At-  
 198 lantic. The rationale for including some of these other regions of the North Atlantic is as  
 199 follows. The southwest corner of the refinement region was chosen to contain the full Gulf  
 200 Stream all the way from the Florida Straits. The southeast corner was chosen to include  
 201 an important region of synoptic eddy wave breaking. The northwest corner was chosen  
 202 to include the entirety of the Labrador Sea and Greenland. The northeast corner was  
 203 chosen to simulate polar lows in the refinement region and to include important regions  
 204 of sea-ice variability, the atmospheric response to which we plan to look at in subsequent  
 205 work.

206 All simulations were performed on the Cheyenne Supercomputer (Computational  
 207 and Information Systems Laboratory, 2019). Based on the known scaling behavior of variable-  
 208 resolution CAM-SE (discussed in Herrington et al., 2022), we chose a relatively small num-  
 209 ber of nodes (30 nodes; 1080 cores) for the NATLx8 simulations for efficiency, because  
 210 we were compute-time rather than throughput limited. The computational cost (includ-  
 211 ing I/O) was approx. 71,000 core-hours per simulated year (CHPSY) for 50-day simu-  
 212 lation segments, which completed in approx. 9 hours and were chosen to be under the  
 213 12-hour wall time. For NATLx4, the computational cost was approx. 21,500 CHPSY for  
 214 6-month simulation segments using 30 nodes, which completed in approx. 10 hours and  
 215 were chosen to be under the 12-hours wall time. For NE30, the computational cost was  
 216 approx. 1,900 CHPSY for 6-month simulation segments using 4 nodes, which completed  
 217 in approx. 7 hours. We thus found that NATLx4 and NATLx8 have  $11\times$  and  $37\times$  in-  
 218 creases in cost compared to NE30, respectively, where this includes I/O and the num-  
 219 ber of nodes used was changed according to what was practical. In total, approximately  
 220 10 million core-hours were used for the simulations in this paper. These simulations also  
 221 serve the purpose of testing this new variable-resolution grid, with additional simulations  
 222 forthcoming.

### 223 2.3 Idealized Specified-SST Experiments

224 For each grid (NATLx8, NATLx4, and NE30) we run a reference simulation with  
 225 year-2000 radiative forcing, year-1850 land-use, and a specified seasonally varying SST  
 226 climatology. The specified climatological SSTs and sea ice concentrations are computed  
 227 from an average over years 1995-2005 in a merged dataset composed of the Hadley Cen-  
 228 ter’s SST/sea-ice version 1.1 and the NOAA Optimal Interpolation analysis version 2  
 229 (Hurrell et al., 2008). These boundary conditions are imposed at  $1^\circ$  spatial resolution  
 230 and monthly time resolution and are interpolated to the atmospheric-model grid and daily  
 231 time resolution by the CESM coupler. All simulations are started from January 1st fol-  
 232 lowing a spin-up procedure needed to generate stable initial conditions (Supporting In-  
 233 formation). Four years of further spin-up are excluded from each simulation due to an  
 234 extended period of stratospheric spin-up in our simulations (Fig. S1 in Supporting In-  
 235 formation). NATLx8 and NATLx4 simulations are extended to February 28th of model  
 236 year 35, accumulating climate statistics over a total of 30 years per simulation. NE30  
 237 simulations are extended to February 28th of model year 55, accumulating climate statis-  
 238 tics over a total of 50 years per simulation.

239 In addition to the reference simulations (referred to as REF throughout the rest  
 240 of the text), we run two SST anomaly experiments for each grid. In the first, we increase  
 241 the SST gradient over the longitudes  $42\text{-}72^\circ\text{W}$  in the Gulf Stream region, with SST anoma-  
 242 lies linearly varying from  $2^\circ\text{C}$  at  $38^\circ\text{N}$  to  $-2^\circ\text{C}$  at  $44^\circ\text{N}$  (Fig. 2a; referred to as GRAD  
 243 throughout the rest of the text). In the second, SSTs are raised by  $2^\circ\text{C}$  everywhere within  
 244 the Gulf Stream box ( $42\text{-}72^\circ\text{W}$ ,  $38\text{-}44^\circ\text{N}$ ) (Fig. 2b; referred to as WARM throughout the  
 245 rest of the text). In both cases, the SST anomalies are imposed in all seasons on top of  
 246 the seasonally varying climatology described in the previous paragraph. The spatial ex-



**Figure 2.** SST anomalies (shading) imposed in all months in the two idealized experiments: (a) The SST gradient anomaly experiment (GRAD); (b) the warm SST anomaly experiment (WARM). The DJF-mean SST climatology is shown in contours, with a contour interval of  $1^{\circ}\text{C}$ .

247 tent of the imposed SST anomalies was chosen based on the large SST variance observed  
 248 in this region (S. M. Wills et al., 2016).

249 The motivation for GRAD was to increase the SST gradient across the Gulf Stream.  
 250 However, when we found that the results did not fit with our expectations for increased  
 251 baroclinicity, we ran the WARM experiments to test whether the simulated response re-  
 252 sulted from the increase in SST gradient or simply from the warming of SSTs in the south-  
 253 ern part of the Gulf Stream region. Our results will show that for the NATLx8 and NATLx4  
 254 grids the WARM simulations produce surprisingly similar results to those in the GRAD  
 255 experiments, suggesting that the warm SSTs in the southern part of the Gulf Stream re-  
 256 gion are the most important aspect of the imposed SST anomalies. Many other stud-  
 257 ies have used a smoothing of SSTs to reduce the SST gradient across the Gulf Stream  
 258 (Nakamura et al., 2008; Kuwano-Yoshida et al., 2010; Parfitt et al., 2016; O’Reilly et al.,  
 259 2016; O’Reilly et al., 2017; Sheldon et al., 2017; Vanni ere et al., 2017; Tsopouridis et al.,  
 260 2021) without introducing abrupt SST jumps, such as occur at  $38^{\circ}\text{N}$  and  $44^{\circ}\text{N}$  in our  
 261 simulations. In hindsight, we believe that this type of SST anomaly experiment may be  
 262 easier to interpret than the ones used here. Nevertheless, the results of our idealized SST  
 263 anomaly experiments (see Section 3) already provide substantial insight into how the at-  
 264 mospheric response to midlatitude SSTs varies with resolution.

265 Output is saved at monthly, daily, and 6-hourly temporal resolution. All output  
 266 is conservatively remapped to a common  $1.25^{\circ}$  longitude  $\times$   $0.94^{\circ}$  latitude grid (referred  
 267 to as f09) for plotting; the f09 grid has a grid spacing of 100-110 km (i.e., comparable  
 268 to NE30) in the midlatitudes. In Section 3.5, we also utilize conservative remapping to  
 269 a  $2.5^{\circ}$  longitude  $\times$   $1.9^{\circ}$  latitude grid (referred to as f19) to separate between large-scale  
 270 and mesoscale anomalies. Unless otherwise indicated, 3D output is linearly interpolated  
 271 from the model’s hybrid coordinates to pressure coordinates (with 31 pressure levels) for  
 272 plotting.

273 Our NATLx8 simulations exhibit large excursions in the global-mean stratospheric  
 274 temperature, both at the beginning of the simulation and following model stability prob-  
 275 lems in model-years 10 and 11 of NATLx8-WARM and NATLx8-REF, respectively (Sup-  
 276 porting Information Fig. S1a). There are corresponding anomalies in the stratospheric  
 277 polar vortex strength, which stand out more from the background variability more in sum-  
 278 mer than in winter (Supporting Information Fig. S1). These excursions appear to be caused  
 279 by reductions in the dynamics timestep that were made to keep the model stable dur-  
 280 ing spin-up and for 1 to 2 days during the aforementioned model stability problems, but  
 281 they persist for several years after the timestep has been returned to its default value.  
 282 Because the stratospheric anomalies in the first 4 years affect all NATLx8 simulations,

283 we discard these years as spinup from the rest of our analysis. Only NATLx8-WARM  
 284 is affected by large stratospheric anomalies in years 10-16, so in this case we simply test  
 285 the sensitivity of our key result to the exclusion of the 6 affected DJFs, finding that it  
 286 is unaffected by the exclusion of this period (Supporting Information Fig. S2). We there-  
 287 fore show averages that include this period in all figures in the main text. More infor-  
 288 mation about these stratospheric excursions is provided in the Supporting Information.

### 289 3 Results

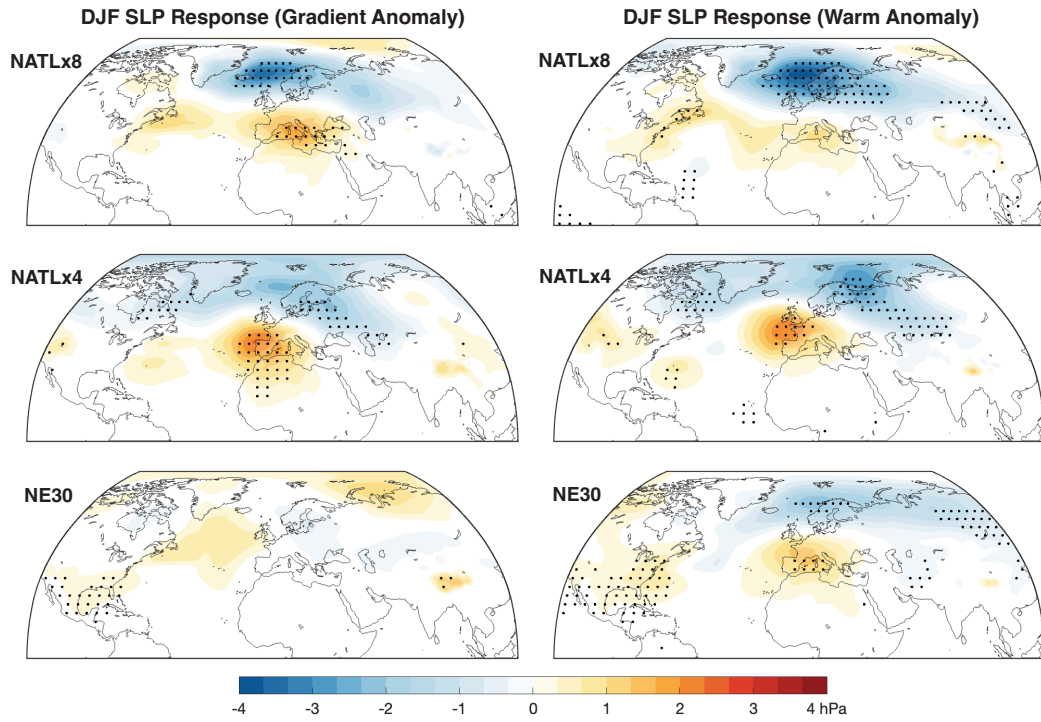
290 Motivated by the potential implications for seasonal-to-decadal predictability, we  
 291 focus our analysis on the response to the imposed SST anomalies, discussing aspects of  
 292 how the climatology changes with resolution when it is relevant. The NAO in December-  
 293 January-February (DJF) is a particularly important target for predictions, and as we will  
 294 show, it has a large response to the SST forcing in the NATLx8 and NATLx4 simula-  
 295 tions that is weaker or absent in the NE30 simulations. We therefore focus our analy-  
 296 sis on DJF (See Supporting Information Fig. S3 for a look at the NAO response in other  
 297 seasons). We note that in contrast to DJF, SON shows a larger NAO response in NE30  
 298 than in the higher resolution simulations, but we leave further investigation of this anomaly  
 299 for future work.

#### 300 3.1 Large-Scale Circulation Response

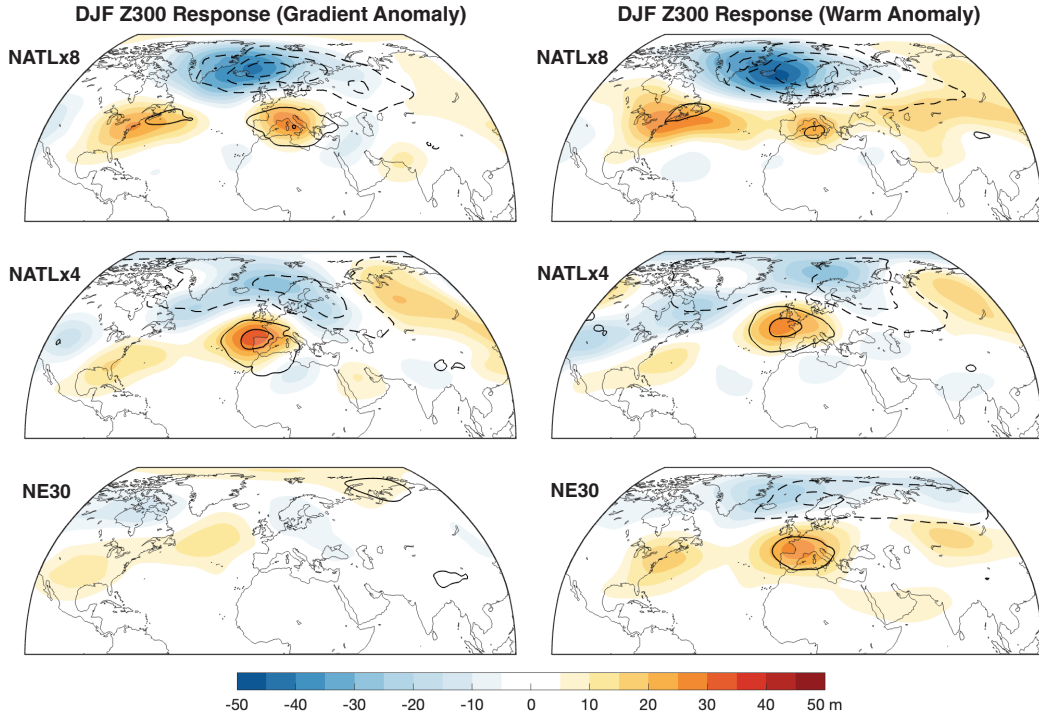
301 To visualize the large-scale circulation response to North Atlantic SST forcing in  
 302 winter (DJF), we first show the DJF sea-level pressure response (Fig. 3). In the high-  
 303 est resolution (14-km) NATLx8 simulations, there is a large East-Atlantic-intensified NAO-  
 304 like response to the SST anomalies in both the GRAD and WARM experiments. It in-  
 305 cludes a large ( $\sim 4$  hPa) negative SLP anomaly centered in the Norwegian Sea and a  
 306 weaker positive SLP anomaly with lobes over the Gulf Stream and Mediterranean. The  
 307 SLP response to the warm SST anomaly (WARM) has a similar spatial pattern in the  
 308 (111-km) NE30 simulations but is weaker in magnitude, especially in the Norwegian Sea.  
 309 The response to the SST gradient anomaly (GRAD) is very weak in NE30, with a com-  
 310 pletely different spatial pattern. If the NATLx8 responses can be thought of as the cor-  
 311 rect response, then the SLP responses in the (28-km) NATLx4 simulations represent an  
 312 improvement compared to NE30, but they show a different spatial pattern with much  
 313 weaker negative anomalies in the Norwegian Sea and stronger positive anomalies over  
 314 Western Europe. Note that the magnitudes of the SLP responses and other key circu-  
 315 lation responses will be tabulated in Section 4.

316 To test the significance of these responses with respect to internal variability, we  
 317 recompute differences from bootstrapped resampling of the three simulations (REF, GRAD,  
 318 and WARM) at each resolution. Differences are computed between averages of  $n' = n(1 -$   
 319  $a)/(1+a)$  resampled years, where  $n$  is the number of years used to compute the response  
 320 (i.e., 30 for NATLx8/NATLx4 and 50 for NE30) and  $a$  is the absolute value of the zonal-  
 321 mean 1-year autocorrelation of seasonal averages at each latitude. The autocorrelation  
 322 factor corrects for the presence of autocorrelation in the original averages that is not present  
 323 in the resampled averages (Trenberth, 1984). We find that a large region of negative SLP  
 324 anomalies in the Norwegian Seas is significant (0.1 significance level based on two-tailed  
 325 t-test; stippling in Fig. 3) in both NATLx8 simulations. The positive SLP anomaly in  
 326 the Mediterranean is also significant in NATLx8-GRAD. NATLx4 shows similar regions  
 327 of significant SLP anomalies (Western Europe and Scandinavia/western Russia) in both  
 328 simulations. NE30-WARM shows only small regions of significant SLP anomalies over  
 329 the North Atlantic and Europe, even with its longer 50-year averages, however, both NE30  
 330 simulations show a large region of weakly positive but significant SLP anomalies over  
 331 the southeast U.S. These results are similar if 30-year averages are used instead of 50-  
 332 year averages for NE30 (Fig. S4 in Supporting Information).





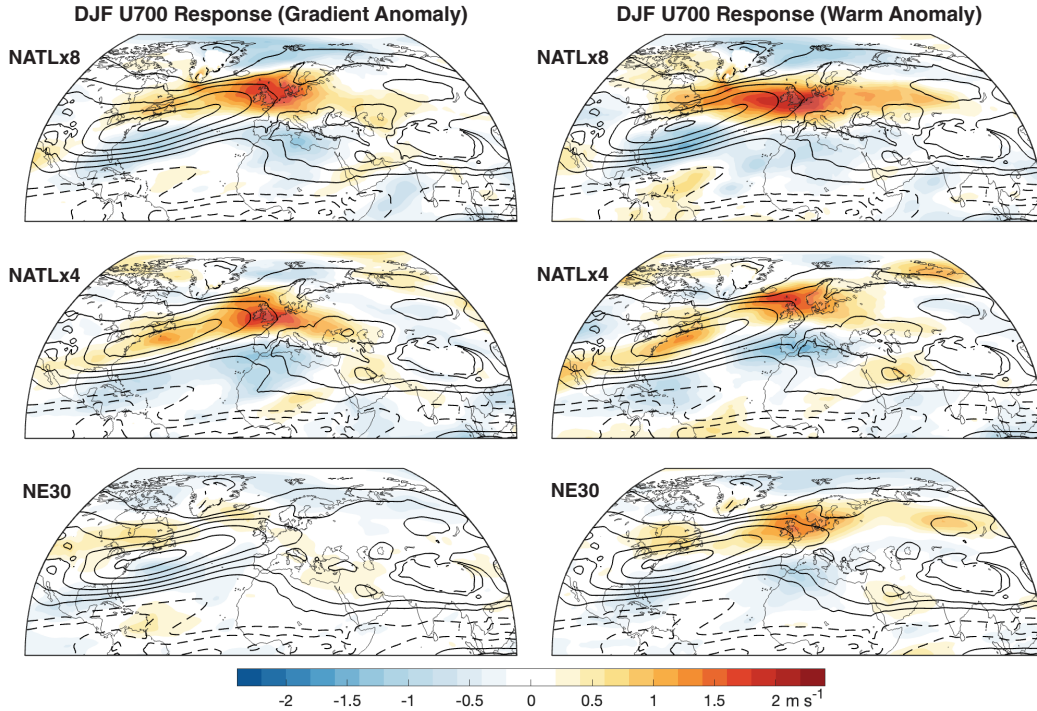
**Figure 3.** DJF Sea-level pressure (SLP) response to an SST gradient anomaly (GRAD-REF; left) and a warm SST anomaly (WARM-REF; right) in the Gulf Stream, in 3 different configurations of CAM-SE: (top) NATLx8, with 14-km resolution in the North Atlantic, (middle) NATLx4, with 28-km resolution in the North Atlantic, and (bottom) NE30, with global 111-km resolution. Anomalies are the difference of 30-year averages in NATL and 50-year averages in NE30. Stippling denotes anomalies that are significant (0.1 significance level) compared to internal variability, diagnosed by bootstrap sampling an equivalent number of independent seasonal averages, accounting for the autocorrelation between seasonal averages as described in the text, and then applying a two-tailed t-test.



**Figure 4.** DJF 300 hPa geopotential height (Z300) response (shading) to an SST gradient anomaly (GRAD-REF; left) and a warm SST anomaly (WARM-REF; right) in the Gulf Stream, in 3 different configurations of CAM-SE: NATLx8, with 14-km resolution in the North Atlantic, NATLx4, with 28-km resolution in the North Atlantic, and NE30, with global 111-km resolution. Anomalies are the difference of 30-year averages in NATL and 50-year averages in NE30. Black contours show the DJF SLP response, as shown in Fig. 3, with a contour interval of 1 hPa; negative anomalies are dashed and the zero contour is omitted.

333 Notably, the similar spatial patterns of SLP response between NE30-WARM and  
 334 NATLx8-WARM, but with much larger magnitudes in NATLx8-WARM, is exactly what  
 335 we should see for this to help alleviate the signal-to-noise paradox (Eade et al., 2014; Scaife  
 336 & Smith, 2018; D. M. Smith et al., 2020). The signal-to-noise paradox is based on the  
 337 finding that models predict the observations well for some quantities (e.g., the NAO),  
 338 but with a reduced amplitude of anomalies, such that the scaled-up ensemble mean has  
 339 more skill in predicting the observations than would be expected from the skill in pre-  
 340 dicting individual ensemble members. Our results suggest that increasing the resolution  
 341 of atmospheric models to better resolve frontal processes could increase the magnitude  
 342 of responses to SST anomalies. In modeling configurations that skillfully predict SSTs,  
 343 the increase in resolution would also increase the magnitude of predictable SLP anom-  
 344 alies, reducing the signal-to-noise paradox. Our results indicate that  $1/4^\circ$  spatial resolu-  
 345 tion may not be enough to recover the full strength of the atmospheric response to mid-  
 346 latitude SST anomalies, and we have no way of knowing whether the response has con-  
 347 verged at  $1/8^\circ$  spatial resolution. The implications for the signal-to-noise paradox will  
 348 be discussed further in Section 4.

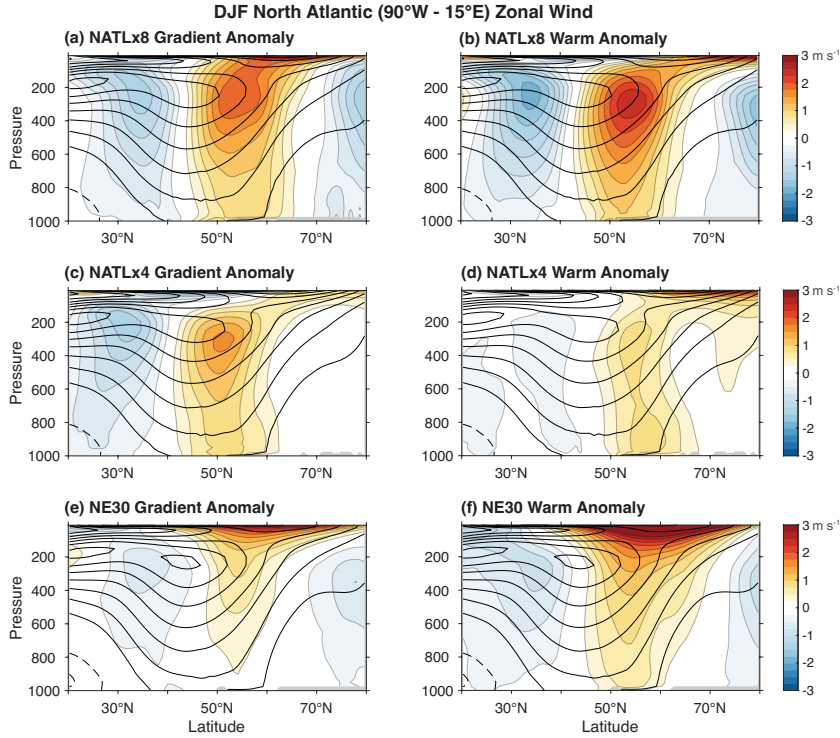
349 The difference in circulation response between NATLx8 and NATLx4 is even more  
 350 apparent in the upper troposphere, as seen in the 300-hPa geopotential height (Z300)  
 351 responses (Fig. 4). The NATLx8 Z300 responses show similar spatial patterns to the SLP  
 352 response, with a westward phase shift indicating an upward propagating stationary wave.



**Figure 5.** DJF 700 hPa zonal wind (U700) response (shading) to an SST gradient anomaly (GRAD-REF; left) and a warm SST anomaly (WARM-REF; right) in the Gulf Stream, in 3 different configurations of CAM-SE: NATLx8, with 14-km resolution in the North Atlantic, NATLx4, with 28-km resolution in the North Atlantic, and NE30, with global 111-km resolution. Black contours show the climatology in the reference simulation (REF) with a contour interval of  $3 \text{ m s}^{-1}$ ; negative anomalies are dashed and the zero contour is omitted). Anomalies are the difference of 30-year averages in NATL and 50-year averages in NE30.

353 The NATLx4 Z300 response shows weaker anomalies with no phase shift compared to  
 354 the SLP response, indicating a stationary wave that is decaying with height. In NE30,  
 355 there are strong westward shifted anomalies in the WARM experiment but weak anom-  
 356 alies with no phase shift in the GRAD experiment.

357 Next, motivated by the finding that models have much weaker decadal variability  
 358 in the zonal-wind at 700 hPa (U700) than is found in reanalysis (Simpson et al., 2018),  
 359 we investigate the U700 response to SST anomalies at each resolution (Fig. 5). All sim-  
 360 ulations except NE30-GRAD show a stronger eastward extension of the climatological  
 361 winds into the UK and Scandinavia in response to the SST anomalies. This response is  
 362 strongest in NATLx8-WARM, then has similar magnitudes in NATLx8-GRAD, NATLx4-  
 363 WARM, and NATLx4-GRAD, but with the largest area of strong anomalies in NATLx8-  
 364 GRAD. NE30-WARM shows a spatially similar but weaker response in this region. These  
 365 results indicate that multi-decadal variance in U700 could increase at higher resolution  
 366 in cases where SSTs vary on multi-decadal timescales. The U700 response varies more  
 367 with resolution in the Gulf Stream SST forcing region: NATLx8 and NE30 show a pole-  
 368 ward wind shift in this region that is stronger in NATLx8 than in NE30, however, NATLx4  
 369 instead shows an intensification of the zonal winds near their climatological maximum.  
 370 These differences don't appear to stem from differences in the climatological winds, which  
 371 are similar across the different resolutions (black contours in Fig. 5).

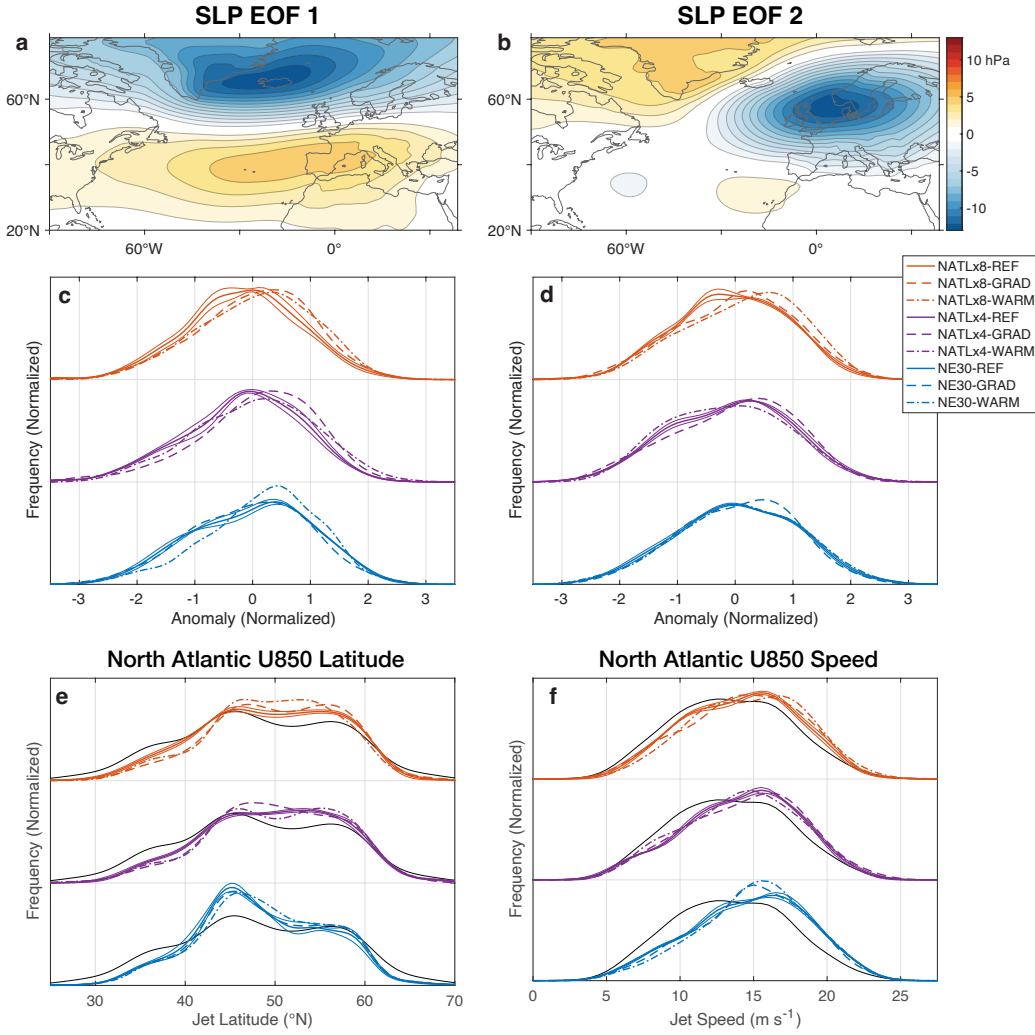


**Figure 6.** Same as Fig. 5, but for the DJF zonal-mean zonal winds over the Atlantic sector ( $90^{\circ}\text{W}-15^{\circ}\text{E}$ ) as a function of latitude and pressure. The contour interval for the climatology is  $4\text{ m s}^{-1}$ ; negative anomalies are dashed and the zero contour is omitted.

372 Due to the increasing amplitude of anomalies with height in NATLx8 and the  
 373 cancellation of anomalies between the eastern and western North Atlantic in NATLx4, the  
 374 zonal-mean zonal winds over the Atlantic sector show much larger anomalies in response  
 375 to SST forcing in NATLx8 than in any of the the lower resolution simulations (Fig. 6).  
 376 All simulations show a poleward shift of the North Atlantic jet, but with different mag-  
 377 nitudes and vertical structures. There are some minor differences in the climatology of  
 378 the North Atlantic zonal winds with resolution, most notably stronger maximum winds  
 379 in the eddy-driven jet in NE30 compared to NATLx4 and NATLx8 and stronger winds  
 380 in the “neck region” (i.e., at  $\sim 100\text{ hPa}$  between the eddy-driven jet and the stratospheric  
 381 polar vortex) in NATLx4 compared to NE30 and NATLx8 (black contours in Fig. 6).

### 382 3.2 Projection onto modes of internal variability

383 To characterize how the large-scale circulation response to SST anomalies projects  
 384 onto the dominant modes of variability, we compute the EOFs of pentadal (5-day-mean)  
 385 SLP in the North Atlantic ( $90^{\circ}\text{W}-40^{\circ}\text{E}$ ,  $15-80^{\circ}\text{N}$ ). We compute the EOFs using 29-years  
 386 (due to missing daily data in one year) of DJF data from each of the 9 simulations to  
 387 obtain a common set of EOFs that explain the variability across all simulations (the EOFs  
 388 computed separately for each resolution are shown in Supporting Information Fig. S5).  
 389 The leading EOF (24% variance explained) represents the NAO (Fig. 7a). The second  
 390 EOF (18% variance explained) shows a low pressure anomaly centered in the North Sea  
 391 and is similar to the East Atlantic pattern (Fig. 7b). The magnitude of both patterns  
 392 is between 12 and 13 hPa, already giving a sense that the  $\sim 4\text{ hPa}$  time-mean anom-  
 393 alies in response to SST anomalies are not small, even compared to synoptic (pentadal)  
 394 variability.



**Figure 7.** (a)-(d) Empirical orthogonal functions (EOFs) of pentadal-mean sea-level pressure (SLP) anomalies in DJF across all 9 simulations, where anomalies are with respect to the average climatology over all 9 simulations and thus include climatological differences. (a) EOF 1, (b) EOF2, (c) probability distribution of principal component 1 in each simulation, (d) probability distribution of principal component 2 in each simulation. The EOFs shown in (a) and (b) are equivalent to the anomaly when the associated principal component is equal to 1. (e) Normalized probability distributions of the pentadal-mean latitude of maximum North Atlantic jet speed during DJF in each simulation and (f) the same for the jet speed at this maximum. The North Atlantic jet is defined as the zonal-mean of the zonal wind at 850 hPa over 0-60°W. In (e) and (f), the black lines show the same analysis applied to ERA5 over 1979-2022. Probability distributions are estimated with kernel density estimation (Botev et al., 2010). Sampling uncertainty in the probability distributions is estimated by splitting each simulation into three segments and dividing the variance in the probability distribution across the segments by 3; the resulting 1-standard-deviation spread is shown for the REF simulations as thin solid lines.

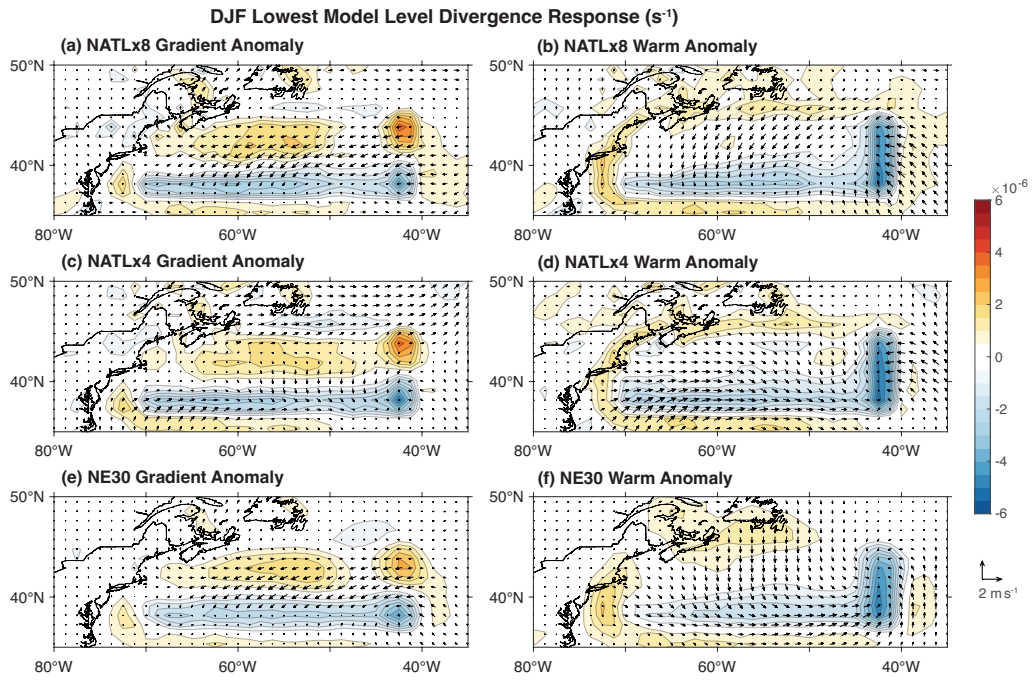
395 The distribution of principal components are shown separately for each simulation  
 396 in Figs. 7c and 7d. In NATLx8, there is a positive shift in the mean of 0.19 (0.21) stan-  
 397 dard deviations in EOF 1 and of 0.07 (0.20) standard deviations in EOF2 in response  
 398 to SST anomalies, for GRAD (WARM) compared to REF. These are responses compared  
 399 to the pentadal variability, and when compared to interannual variability they instead  
 400 correspond to shifts in the mean of 0.41 (0.44) standard deviations in EOF 1 and of 0.21  
 401 (0.61) standard deviations in EOF2, for GRAD (WARM) compared to REF. NATLx4-  
 402 GRAD (NATLx4-WARM) show a similar positive shift in the mean of EOF1 of 0.52 (0.42)  
 403 interannual standard deviations but a smaller shift in the mean of EOF 2 of 0.15 (-0.15)  
 404 interannual standard deviations. In NE30-WARM, the probability of negative EOF 1  
 405 values is reduced in favor of an increase in the probability of weakly positive EOF 1  
 406 values, near the peak of the distribution, corresponding to a shift in the mean of 0.35 in-  
 407 terannual standard deviations, but it has no meaningful change in the distribution for  
 408 EOF 2 (a mean shift of 0.14 interannual standard deviations). NE30-GRAD does not  
 409 show much of a shift in either EOF, with mean shifts of -0.12 and 0.13 interannual stan-  
 410 dard deviations in EOFs 1 and 2, respectively. Overall, this analysis shows that the SST  
 411 anomalies both lead to large ( $\sim 0.5$  interannual-standard-deviation) anomalies in the lead-  
 412 ing mode of SLP variability (i.e., the NAO) in NATLx8 and NATLx4 that are weaker  
 413 or absent in NE30, and that NATLx8-WARM shows a much larger response in EOF 2  
 414 than any of the other simulations.

415 North Atlantic circulation variability has also been characterized by the latitude  
 416 of the jet maximum, which has been shown to exhibit regime-like behavior not appar-  
 417 ent from the EOFs of SLP (Woollings et al., 2010; White et al., 2019; Strommen et al.,  
 418 2019; Strommen, 2020; Dorrington et al., 2022). Following Strommen (2020), we com-  
 419 pute the North Atlantic jet latitude as the latitude of the maximum in the zonal-mean  
 420 850-hPa zonal winds in the North Atlantic (0-60°W). We use pentadal averages in place  
 421 of the 9-day running mean used in Strommen (2020). NATLx8 has the most realistic struc-  
 422 ture of the jet latitude probability distribution compared to ERA5 Reanalysis (Hersbach  
 423 et al., 2020) (Fig. 7e), but all 3 resolutions of CAM6-SE show too little occurrence of  
 424 the southernmost jet latitude peak at 35°N. The 45°N jet latitude peak is too strong in  
 425 NE30, whereas it is more realistic in NATLx4 and NATLx8. Both NATLx4 and NATLx8  
 426 have a relatively larger probability (compared to NE30 and ERA5) of jets occurring at  
 427 the northern peak, the presence of which has been linked to Greenland topography and  
 428 Greenland tip-jet events (White et al., 2019). Overall, there is some indication that the  
 429 regime-like behavior of jet latitude increases with resolution (cf. Strommen, 2020), which  
 430 is apparent in the less peaked probability distributions in NATLx4 and NATLx8 com-  
 431 pared to NE30. In terms of jet speed, NATLx8 is again most realistic compared to ERA5  
 432 reanalysis (Fig. 7f).

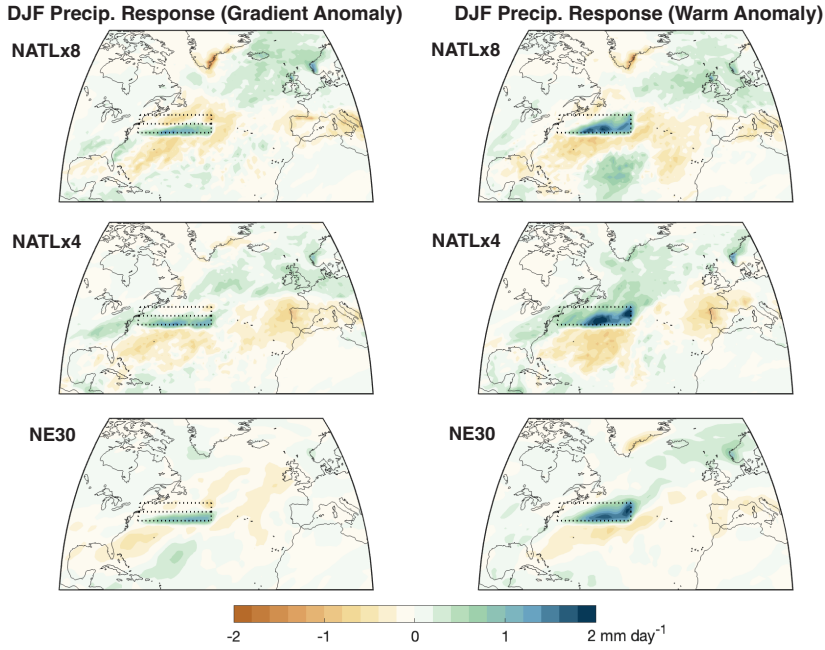
433 In response to both SST anomalies, NATLx8 and NATLx4 show increases in the  
 434 probability of jets at the midlatitude and northern peaks at the expense of jets at the  
 435 southern peak (Fig. 7e) and a slight shift towards stronger jet speeds (Fig. 7f). In con-  
 436 trast, NE30-WARM (and to a lesser extent NE30-GRAD) shows a more peaked jet speed  
 437 distribution, a poleward shift of the midlatitude peak, an increase in the probability of  
 438 jets at the northern peak, and no change in the probability of jets at the southern peak.  
 439 Overall, this shows that the circulation response to SST anomalies is more complex than  
 440 a simple mean shift in circulation and it is associated with a shift in probability of oc-  
 441 currence of the underlying circulation regimes.

### 442 3.3 Air-Sea Interactions and Cross-Front Circulation Response

443 As a first step in analyzing the mechanisms for the large NAO-like response to SST  
 444 anomalies and its dependence on resolution, we investigate the air-sea interactions and  
 445 the cross-front circulation response in the SST forcing region.



**Figure 8.** DJF near-surface (lowest model level) zonal and meridional wind (arrows) and divergence (shading) response to an SST gradient anomaly (GRAD-REF; left) and a warm SST anomaly (WARM-REF; right) in the Gulf Stream, in 3 different configurations of CAM-SE: NATLx8, with 14-km resolution in the North Atlantic, NATLx4, with 28-km resolution in the North Atlantic, and NE30, with global 111-km resolution. Anomalies are the difference of 30-year averages in NATL and 50-year averages in NE30.

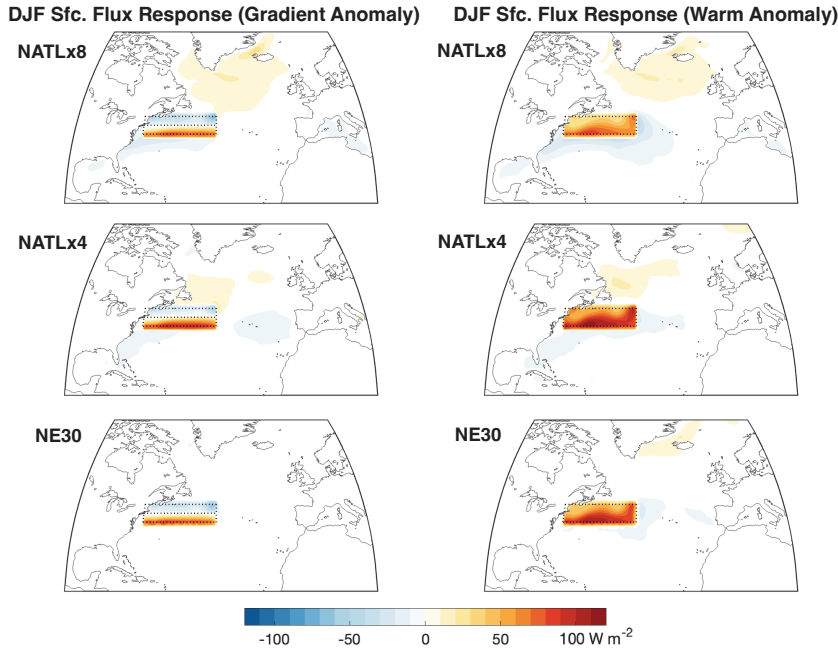


**Figure 9.** DJF precipitation response to an SST gradient anomaly (GRAD-REF; left) and a warm SST anomaly (WARM-REF; right) in the Gulf Stream, in 3 different configurations of CAM-SE: (top) NATLx8, with 14-km resolution in the North Atlantic, (middle) NATLx4, with 28-km resolution in the North Atlantic, and (bottom) NE30, with global 111-km resolution. Anomalies are the difference of 30-year averages in NATL and 50-year averages in NE30.

446 Much of the literature on how ocean resolution impacts the atmospheric response  
 447 to SST anomalies has focused on the near-surface wind divergence (e.g., Small et al., 2014),  
 448 because it is related to the Laplacian of SST through the pressure adjustment mecha-  
 449 nism (Lindzen & Nigam, 1987; Minobe et al., 2008) and to the downwind SST gradient  
 450 by the vertical mixing mechanism (Hayes et al., 1989; Chelton et al., 2001), and because  
 451 both the Laplacian of SST and the downwind SST gradient are sensitive to the ocean  
 452 resolution. However, we find that the near-surface wind divergence response is very simi-  
 453 lar across different atmospheric resolutions (despite differences in the response of the  
 454 individual near-surface wind components; Fig. 8). This suggests that differences in near-  
 455 surface divergence response are not the reason for the differences in large-scale circula-  
 456 tion response with resolution. This is perhaps not surprising considering the strong rela-  
 457 tionship between near-surface divergence and SST, which is kept the same as the at-  
 458 mospheric resolution is varied. Indeed, the spatial pattern of near-surface divergence matches  
 459 well with the downwind SST gradient (leading to large anomalies on the eastern bound-  
 460 ary of the forcing region) and the Laplacian of SST (leading to large anomalies on the  
 461 southern boundary of the forcing region), as expected from these boundary layer theo-  
 462 retical considerations.

463 Precipitation anomalies somewhat resemble the near-surface convergence anomalies  
 464 (Fig. 9), with anomalies over the forcing region of 1-2 mm/day, more than 20% of  
 465 the climatological precipitation in this region. Like the near-surface convergence, they  
 466 do not show large differences across the different resolutions. It therefore does not ap-  
 467 pear that differences in the time-mean precipitation and latent heating are responsible  
 468 for the difference in large-scale circulation response. For example, the experiment with  
 469 the largest precipitation response (NATLx4-WARM) does not have the largest large-scale  
 470 circulation response (cf. Figs. 3-6). Note, however that the precipitation anomalies over





**Figure 10.** Same as Fig. 9, but for DJF surface turbulent (latent + sensible) heat flux.

471 the SST forcing are bounded by dry anomalies to the north in NATLx8, whereas they  
 472 are continuous with enhanced precipitation anomalies to the north in NATLx4. This is  
 473 a qualitative indication that precipitation occurs through local convective process in NATLx8  
 474 versus as part of the larger-scale warm conveyor belt in NATLx4, as will be discussed  
 475 in Section 3.5. Note also that while the time-mean precipitation is similar across resolu-  
 476 tions, the spatial contrast in precipitation and latent heating between different parts  
 477 of cyclones could increase at higher resolution (Schemm, 2023), potentially causing more  
 478 rapid cyclone growth. Downstream of the forcing region, the non-local responses (e.g.,  
 479 in the subpolar North Atlantic and Western Europe) are larger in the NATLx4 and NATLx8  
 480 simulations as a result of the larger large-scale circulation responses, with anomalies in  
 481 the eastern North Atlantic and Europe of up to 10-20% of the climatological DJF pre-  
 482 cipitation in these regions.

483 Given the use of specified-SST experiments, a natural question arises of whether  
 484 the SST anomalies correspond to comparable surface turbulent (latent + sensible) heat-  
 485 flux anomalies as the atmospheric resolution is varied. Similar to what was found for near-  
 486 surface divergence and precipitation, the anomalies in turbulent heat fluxes are differ-  
 487 ent between the GRAD and WARM experiments, but the differences with resolution are  
 488 relatively small (Fig. 10). There is some variation in the magnitude of surface fluxes with  
 489 resolution, especially for the WARM experiment, with the largest values in NATLx4 and  
 490 the smallest in NATLx8. This means that NATLx8 gives the largest large-scale circu-  
 491 lation response despite having the smallest surface heat-flux anomalies. The surface flux  
 492 differences are related to differences in the adjustment of near-surface air temperature,  
 493 with near-surface air temperature anomalies being largest in NATLx8 and smallest in  
 494 NATLx4 (not shown).

495 The differences in air-temperature adjustment over the SST anomalies are also evi-  
 496 dent further into the troposphere; NATLx8-GRAD, NATLx8-WARM, and (to a lesser  
 497 extent) NE30-WARM all show deep warm anomalies over the forcing region ( $42\text{-}72^\circ\text{W}$ ;  
 498 Fig. 11). The differences across the simulations in the magnitude of potential temper-  
 499 ature response over the forcing region mirror the differences in the magnitude of the up-

500 per tropospheric circulation response (cf. Figs. 4 and 6), a simple consequence of ther-  
 501 mal wind balance. Explaining the differences in the free-tropospheric potential temper-  
 502 ature response in the forcing region is therefore key to understanding the differences in  
 503 the large-scale circulation response between simulations. The horizontal spatial struc-  
 504 ture of these deep temperature anomalies can most clearly be seen in Z300 (Fig. 4), which  
 505 is related to the vertically averaged temperature anomaly below 300 hPa. The poten-  
 506 tial temperature responses over the forcing region look different in both NATLx4 exper-  
 507 iments compared to those in the other simulations, with a warm anomaly to the south  
 508 of the forcing region and a cold anomaly to the north (Fig. 11; cf. Fig. 4), consistent  
 509 with the increase in wind speed at the jet maximum that was seen in Fig. 5.

510 Fig. 11 also shows anomalies in the time-mean ageostrophic meridional and verti-  
 511 cal winds over the Gulf Stream SST front. The time-mean upward motion is not very  
 512 different between the different simulations; all experiments show anomalous upward mo-  
 513 tion extending to between 400 and 500 hPa. However, there are large differences in the  
 514 ageostrophic meridional winds. While much of the ageostrophic meridional wind anom-  
 515 alies over the Gulf Stream SST anomalies in NATLx8 appear to make up a closed merid-  
 516 ional circulation, with ascent near 38°N and descent near 45°N, the ascending air anom-  
 517 alies instead turn equatorward in NATLx4 and (to a lesser extent) NE30, similar to what  
 518 was found in Smirnov et al. (2015). Thus only the NATLx8 experiments (and to a lesser  
 519 extent NE30-WARM) have poleward ageostrophic winds in the upper troposphere, which  
 520 can provide an important source of zonal momentum.

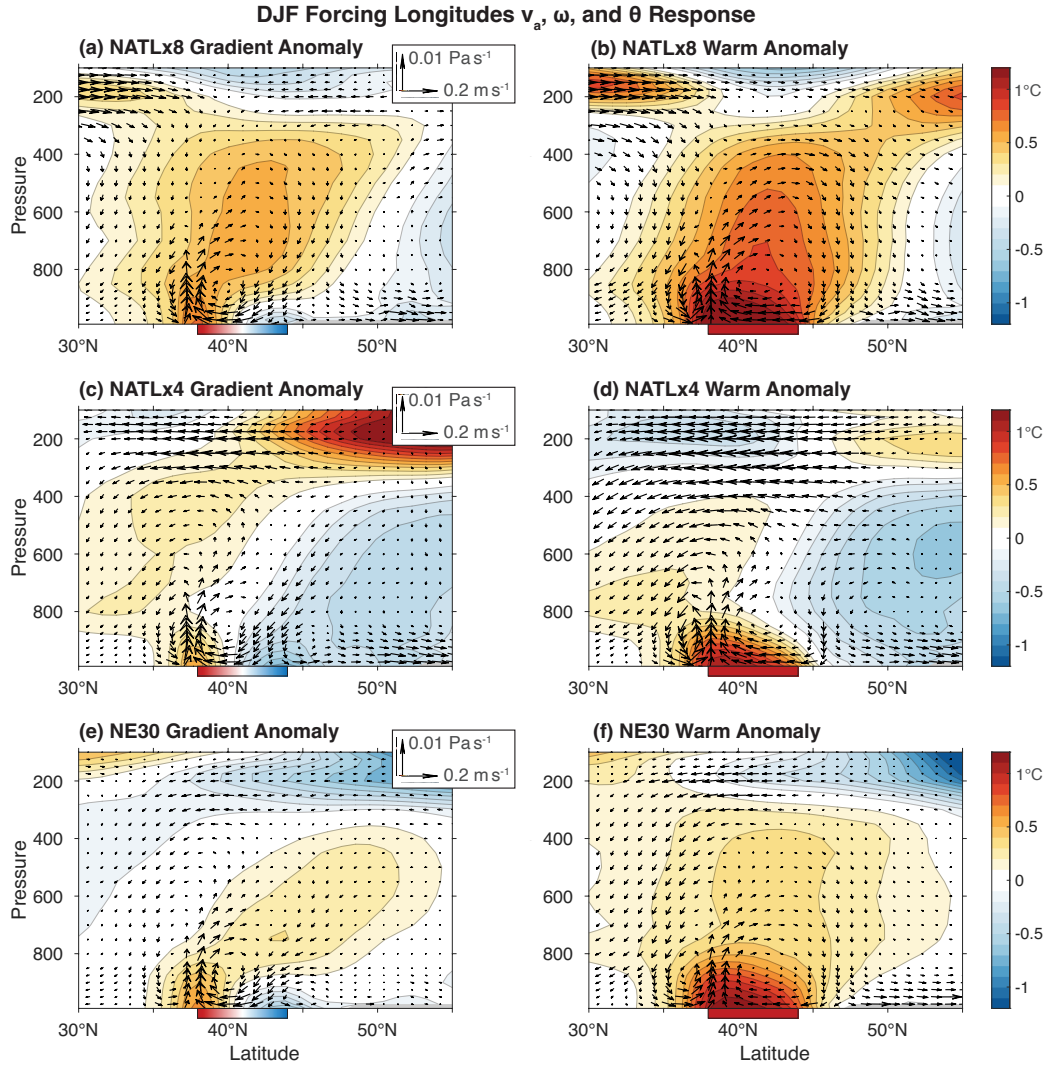
### 521 3.4 Thermodynamic Analysis

522 To gain insight into the maintenance of the deep temperature anomalies in NATLx8-  
 523 GRAD, NATLx8-WARM, and NE30-WARM, we analyze the thermodynamic equation  
 524 for the mid-troposphere (300-800 hPa) over the forcing region (42-72°W; 38-44°N):

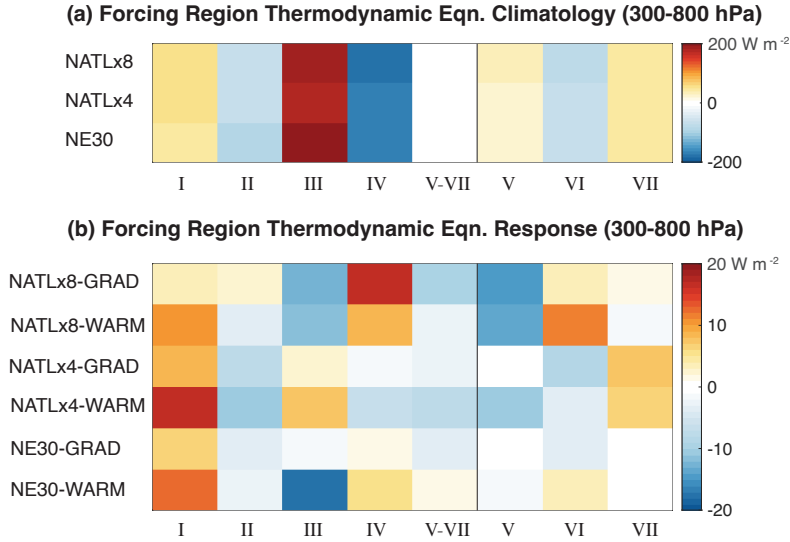
$$\underbrace{\overline{Q}}_{\text{I}} - \underbrace{(\overline{\omega} \partial_p \overline{T} - \kappa \frac{\overline{\omega T'}}{p})}_{\text{II}} - \underbrace{\overline{v} \nabla_y \overline{T}}_{\text{III}} - \underbrace{\overline{u} \nabla_x \overline{T}}_{\text{IV}} - \underbrace{\nabla_x \cdot \overline{u' T'}}_{\text{V}} - \underbrace{\nabla_y \cdot \overline{v' T'}}_{\text{VI}} - \underbrace{(\partial_p (\overline{\omega' T'}) - \kappa \frac{\overline{\omega' T'}}{p})}_{\text{VII}} = 0. \tag{1}$$

525 Here, overbars denote monthly averages, primes denote deviations from the monthly mean,  
 526  $\nabla_x$  and  $\nabla_y$  are the zonal and meridional components of the nabla operator on a sphere,  
 527  $Q$  is the total diabatic heating (including latent heating, radiation, and parameterized  
 528 turbulent diffusion),  $\kappa = R/c_p = 2/7$  is the ratio of the specific gas constant and spe-  
 529 cific heat capacity of dry air, and all other variables follow standard meteorological con-  
 530 ventions. Over the Gulf Stream, the dominant balance is between meridional warm air  
 531 advection (Term III) and zonal advection of cold air off the North American continent  
 532 (Term IV) (Fig. 12a). There is also time-mean upward motion (Term II) and diabatic  
 533 (latent) heating (Term I). The total effect of transient-eddy heat-flux convergence (Terms  
 534 V-VII) is small due to cancellation between heating by zonal and vertical eddy heat trans-  
 535 port and cooling from meridional eddy heat transport. These balances stay roughly the  
 536 same as the resolution is changed.

537 The response of the terms in the thermodynamic equation (in the mid troposphere)  
 538 to the imposed SST anomalies shows more varied behavior across the different resolu-  
 539 tions. All simulations show an increase in latent heating in response to the SST anom-  
 540 alies (Fig. 12b; Term I); this increase in latent heating is largest in the WARM exper-  
 541 iments, owing to a partial compensation by negative anomalies in the northern part of  
 542 the forcing domain in the GRAD experiments (not shown). While the latent heating anom-  
 543 alies are largest in the NATLx4 simulations, matching what was found for precipitation  
 544 and surface fluxes (cf. Figs. 9 and 10), they are compensated in these simulations by larger  
 545 negative anomalies in the vertical advection term (Fig. 12b; Term II). Rather than re-  
 546 sulting from differences in time-mean ascent, which is similar across the resolutions (Fig.  
 547 11), these differences in Term II result from differences in stratification in the ascent re-



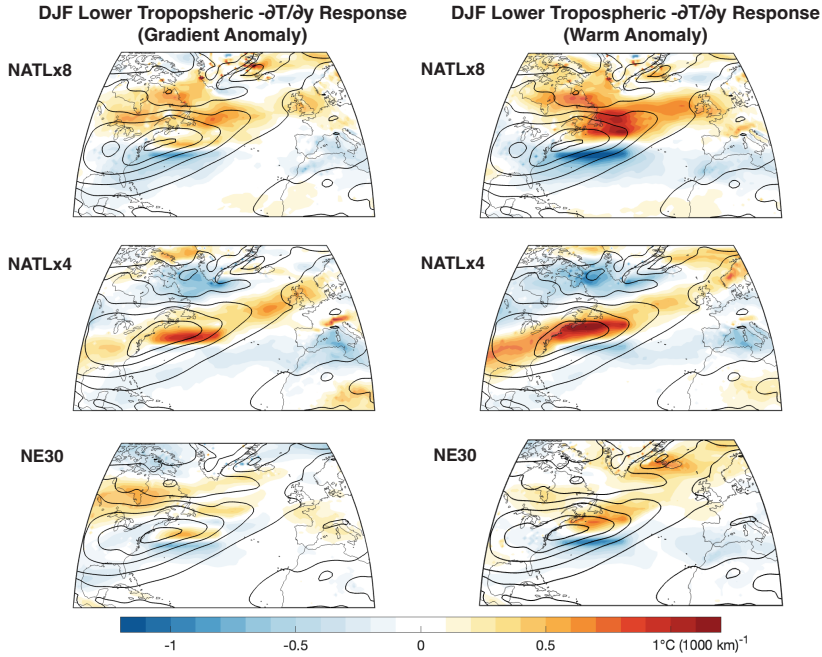
**Figure 11.** Average over the forcing longitudes ( $42\text{--}72^\circ\text{W}$ ) of the DJF potential temperature (shading) and ageostrophic meridional and vertical wind (arrows) response to an SST gradient anomaly (GRAD-REF; left) and a warm SST anomaly (WARM-REF; right) in the Gulf Stream, in 3 different configurations of CAM-SE: (top) NATLx8, with 14-km resolution in the North Atlantic, (middle) NATLx4, with 28-km resolution in the North Atlantic, and (bottom) NE30, with global 111-km resolution. Anomalies are the difference of 30-year averages in NATL and 50-year averages in NE30.



**Figure 12.** Average over the forcing longitudes ( $42\text{--}72^\circ\text{W}$ ) and latitudes ( $38\text{--}44^\circ\text{N}$ ) of (a) the DJF climatology (REF) of the terms in the thermodynamic equation (Eq. 1), including a column showing the sum of terms V-VII, and (b) responses of these terms to an SST gradient anomaly (GRAD-RED) and a warm SST anomaly (WARM-REF) in the Gulf Stream, in 3 different configurations of CAM-SE. Anomalies are the difference of 30-year averages in NATL and 50-year averages in NE30.

548 gion, which decreases in response to the SST anomalies in NATLx8, as well as from anomalous  
 549 time-mean subsidence on the northern and southern edges of the forcing region, which  
 550 is strongest in the NATLx8 SST anomaly experiments. This means that the effective forcing  
 551 from vertical motions after accounting for the cancellation between adiabatic cooling  
 552 and latent heating (Term I + Term II) is similar across different resolutions.

553 Despite broad similarities in the first two terms, the response in the horizontal ad-  
 554 vection terms (Eq. 1; Terms III and IV) are opposite between the simulations with deep  
 555 temperature anomalies (NATLx8 and NE30-WARM) and those with free-tropospheric  
 556 temperature gradient anomalies (NATLx4): NATLx4 shows a strengthening of the cli-  
 557 matological meridional warm-air advection and zonal advection of cold air off the con-  
 558 tinent, whereas NATLx8 and NE30-WARM show a weakening of the climatology (Fig.  
 559 12b; Terms III and IV). The negative meridional advection anomalies (Term III) for NATLx8  
 560 and NE30-WARM result from a combination of northerly wind anomalies (not shown  
 561 in Fig. 11 because they are geostrophic) and weakened meridional temperature gradi-  
 562 ent, whereas the positive anomalies in NATLx4 result primarily from the strengthened  
 563 meridional temperature gradient (Fig. 13). The meridional temperature gradient response  
 564 (Fig. 13) shows a poleward shift in NATLx8 and NE30-WARM but a strengthening near  
 565 its maximum for NATLx4, as was seen for the U700 response (Fig. 5). Similarly, the changes  
 566 in zonal advection (Term IV) in NATLx8 can be partially understood in terms of changes  
 567 in horizontal temperature gradients, with a tropospheric warming over the U.S. eastern  
 568 seaboard reducing the zonal temperature gradient in NATLx8 and NE30-WARM but  
 569 a cooling over Atlantic Canada increasing the zonal temperature gradient in NATLx4  
 570 (Fig. 4). Interestingly, Famooss Paolini et al. (2022) also see switches in sign of the time-  
 571 mean meridional and zonal advection terms between 100-km- and 50-km-resolution mod-  
 572 els, in agreement with the changes between NE30 and NATLx4; however, we see another  
 573 switch in sign of these terms going from NATLx4 (28 km) to NATLx8 (14 km).

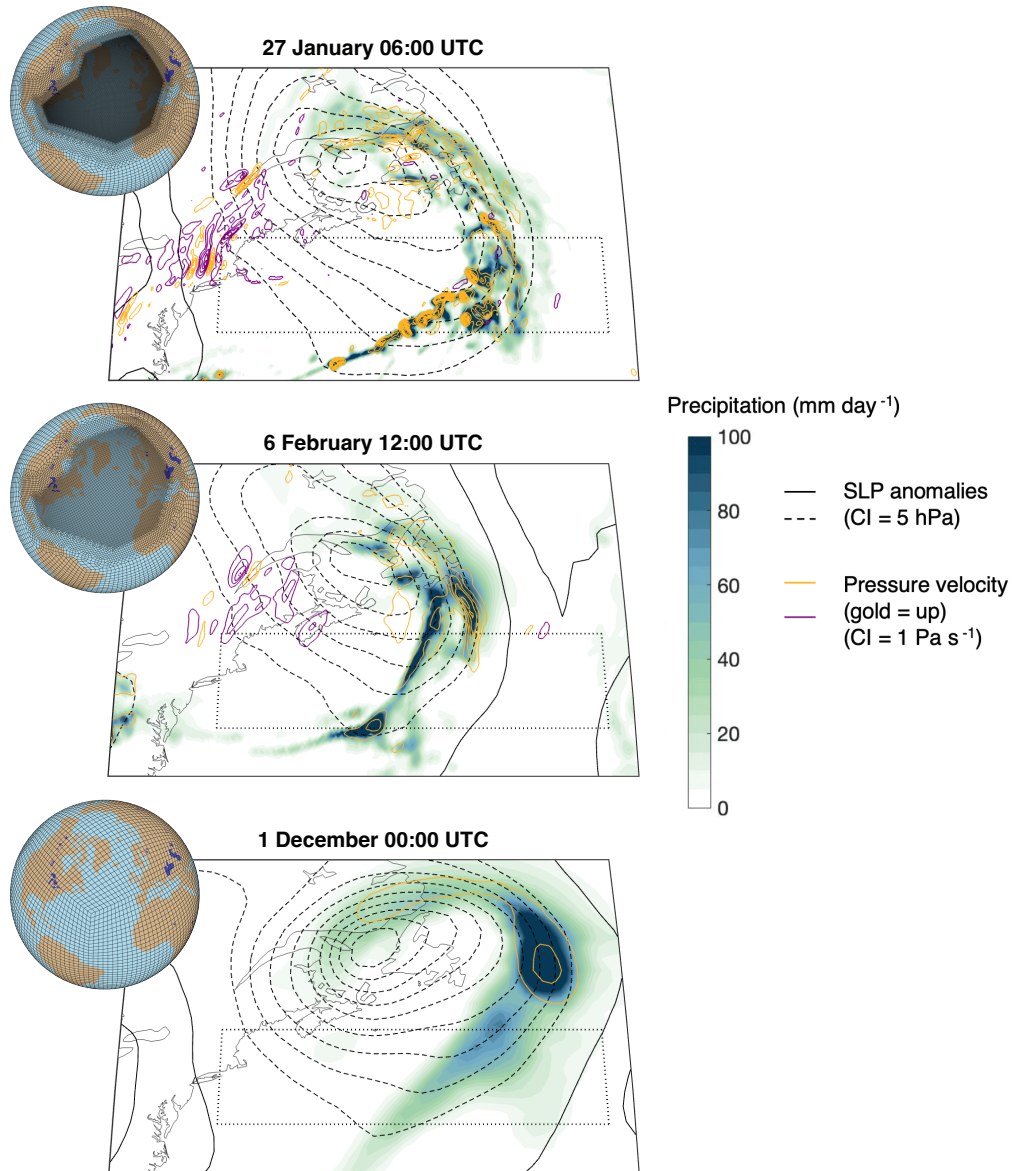


**Figure 13.** Same as Fig. 5, but for the DJF vertically averaged meridional temperature gradient below 500 hPa, with flipped sign such that a poleward decrease in temperature is positive. The contour interval for the climatology (which has been spatially smoothed with a  $1.5^\circ$  Gaussian filter) is  $2^\circ\text{C} (1000 \text{ km})^{-1}$ .

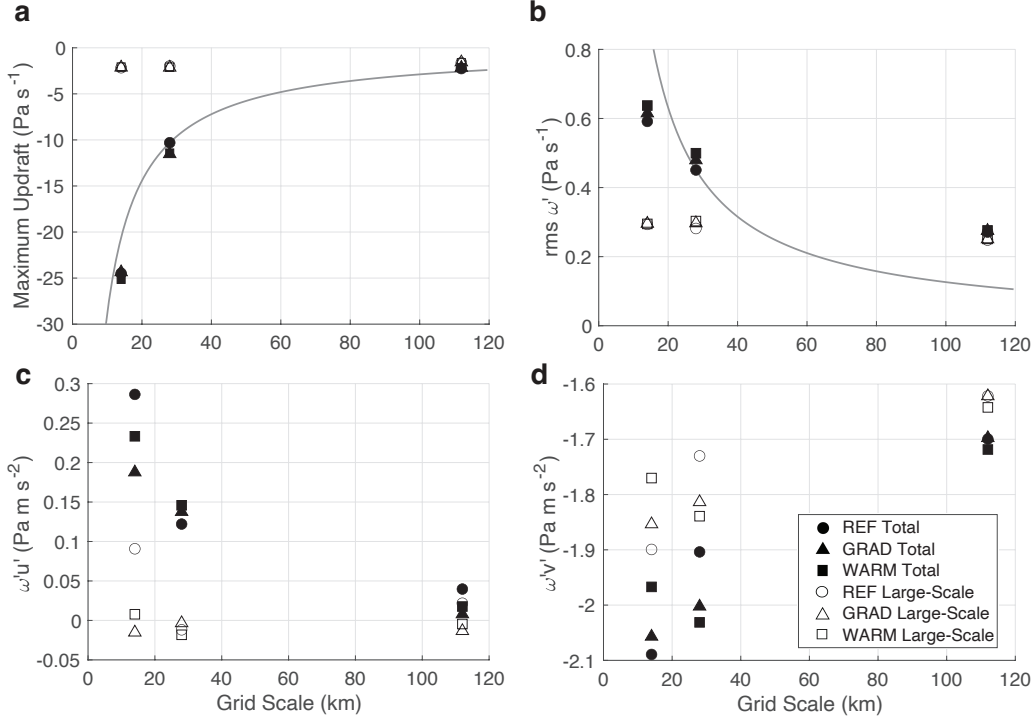
574 Thus far, our analysis of the thermodynamic equation has illustrated differences  
 575 across resolutions in the response of the dominant terms, but it has not provided a defini-  
 576 tive answer to what is driving the deeper warm anomalies in NATLx8 and NE30-WARM.  
 577 This is in part inherent to any analysis of the thermodynamic equation, where individ-  
 578 ual terms influence but are also influenced by the distribution of temperature anoma-  
 579 lies. However, there are only a few terms with anomalously positive tendencies in response  
 580 to SST anomalies in NATLx8 than NATLx4, such that they could explain a larger free-  
 581 tropospheric warming in NATLx8: vertical advection (Term II), zonal advection (Term  
 582 IV), and meridional eddy heat-flux (EHF) convergence (Term VI). It has already been  
 583 discussed how the zonal and vertical advection anomalies are a consequence of the deep  
 584 temperature anomaly, which reduces the zonal temperature gradient and the lapse rate.  
 585 Therefore, in the next section we turn our attention to the responses of meridional EHF  
 586 and other transient-eddy heat fluxes to SST forcing and how they depend on resolution.  
 587 The basic picture that emerges is that frontal processes move heat vertically in NATLx8,  
 588 creating a deep warm temperature anomaly that reduces the meridional temperature gra-  
 589 dient and thus the meridional EHF, the divergence of which would otherwise act to damp  
 590 the temperature anomaly. In contrast, when eddies move heat vertically in NATLx4, they  
 591 do so as part of the cyclone warm conveyor belt, which also moves this heat poleward  
 592 and out of the forcing region.

### 593 3.5 Modification of Transient-Eddy Fluxes

594 Before diving into a quantitative analysis of changes in transient eddy statistics,  
 595 it is helpful to visualize how the transient eddies look qualitatively different between the  
 596 simulations at different resolutions. We therefore show snapshots of low-pressure systems  
 597 passing over the Gulf Stream SST forcing region in one of the simulations at each res-



**Figure 14.** Snapshots of instantaneous total precipitation rate (shading), sea-level pressure (SLP) anomalies from the climatological mean (black contours; dashed negative), and vertical pressure velocity on the model level with average pressure of 610 hPa (gold = up; purple = down) from the WARM experiment at each resolution. Qualitatively similar snapshots are chosen such that they have an extratropical cyclone centered in the Gulf of St. Lawrence, north of the SST forcing region (thin dotted line), during winter. For plotting, precipitation and vertical velocity are interpolated to a uniform  $1/8^\circ$  grid for NATLx8 and NATLx4 and a uniform  $0.7^\circ$  grid for NE30; SLP is interpolated to the  $1.25^\circ$  longitude  $\times$   $0.94^\circ$  f09 grid for all simulations. The element grid is shown alongside each panel, where each element contains  $3 \times 3$  independent computational grid points.



**Figure 15.** Statistics of vertical winds and vertical momentum fluxes in each simulation during DJF, computed from 6-hourly model output and plotted against the horizontal grid scale of the simulations. Statistics are computed over the southern part of the forcing region (42–72°W; 38–41°N) on the model level with average pressure of 610 hPa. (a) Temporal median of the seasonal maximum (instantaneous) updraft speed over the forcing region, expressed in pressure velocity. (b) Root-mean-square temporal variance of vertical pressure velocity over the forcing region. (c,d) Vertical fluxes of (c) zonal and (d) meridional momentum, i.e., the covariance of pressure velocity anomalies with zonal and meridional wind anomalies. In all panels, open symbols show statistics computed from large-scale fields, after interpolation to the  $2.5^\circ \times 1.9^\circ$  f19 grid, such that variations on scales smaller than  $\sim 200$  km are excluded, whereas solid symbols show the statistics computed on the native grid. Here and in Figs. 16–20, statistics are computed over 30 years in NATLx8 and NATLx4, 50 years in NE30-REF and NE30-WARM, and 48 years in NE30-GRAD (due to missing 6-hourly data). Grey lines in (a) and (b) show the  $W \propto D^{-1}$  scaling, with constants chosen to intersect NATLx4-REF.

598 resolution (Fig. 14). The highest resolution NATLx8 shows precipitation organized in frontal  
 599 bands, and there is a well defined cold front with vertical velocities exceeding  $10 \text{ Pa s}^{-1}$ .  
 600 There are also resolved gravity waves apparent in the vertical velocities in the cold sector  
 601 of the cyclone. NATLx4 shows these same basic features but with muted vertical ve-  
 602 locities, especially in the cold front. In comparison to these higher resolution simulations,  
 603 precipitation and vertical velocity in the lower resolution NE30 simulations look much  
 604 more blobular, without well-defined mesoscale features. This section will quantify how  
 605 the large differences in the magnitude and spatial structure of vertical velocities within  
 606 midlatitude cyclones influence transient eddy statistics and help shape the large-scale  
 607 circulation response.

608 The maximum updraft velocities over the Gulf Stream increase with increased resolu-  
 609 tion according to the  $W \propto D^{-1}$  scaling derived in Jeevanjee and Romps (2016) (black

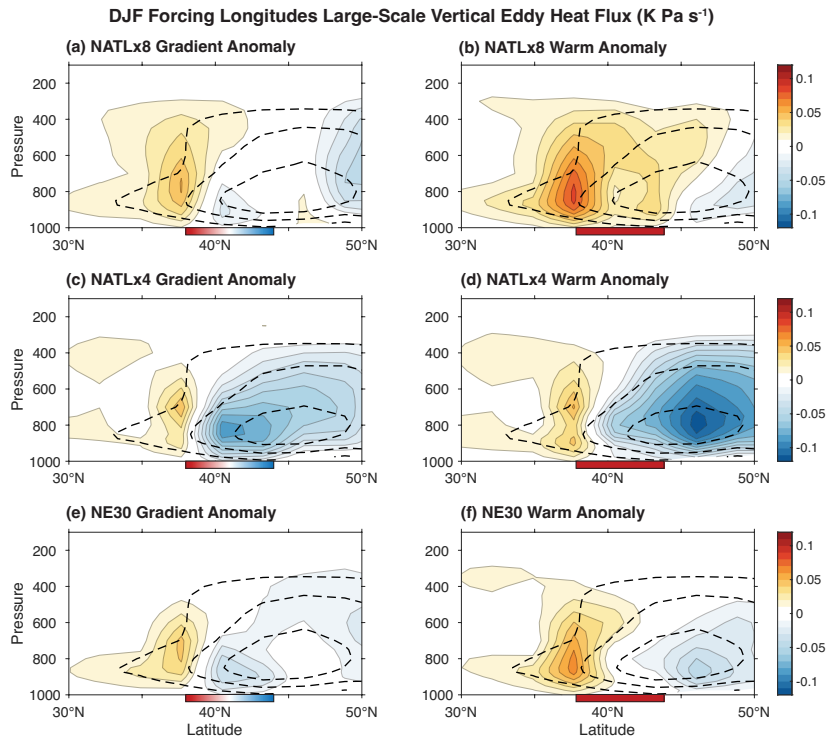
lines in Fig. 15a), where  $W$  is the vertical velocity scale and  $D$  is the horizontal scale of convective updrafts. This is consistent with Herrington and Reed (2018), who showed that this scaling applies across different resolutions of CESM. The reason for this scaling is that buoyancy anomalies develop on smaller scales as the grid scale is reduced and this means that an equivalent buoyancy anomaly will be resisted by a narrower column of air. We actually find that the increase in updraft velocities in our simulations slightly exceeds this scaling (Fig. 15a). As was apparent in Fig. 14, these updrafts occur in mesoscale fronts, and there is therefore little change in the magnitude of large-scale updrafts (open symbols in Fig. 15a). Here, we compute large-scale statistics based on model output that has been conservatively remapped to the  $\sim 200$ -km f19 grid, whereas the full-field statistics (filled symbols in Fig. 14) are computed on the native grid. The vertical velocity variance increases more slowly with resolution than the maximum updraft velocity (Fig. 15b; cf. Fig. 15a), because the area in which the strongest updrafts are occurring reduces with increased resolution. The large-scale vertical velocity variance does increase between NE30 and NATLx4, but most of the vertical velocity variance changes come from scales smaller than 200 km.

In the following discussion of changes in transient eddy fluxes, it is worth bearing in mind that transient eddies include not only synoptic motions and low-frequency variability, as is normally the case in analysis of GCM output, but they also include mesoscale motions such as slantwise convection. Studies based on reanalysis have found evidence that slantwise convection occurs over the Gulf Stream, especially in winter (Korty & Schneider, 2007; Czaja & Blunt, 2011; Sheldon & Czaja, 2014). To quantify the presence of mesoscale shear instabilities such as conditional symmetric instability in our simulations, we examine the vertical momentum fluxes by mesoscale eddies (less than 200 km scales, cf. Sheldon et al. (2017)). The vertical flux of zonal momentum by mesoscale motions (difference between open symbols and closed symbols in Fig. 15c) is positive (downwards) and increases strongly with increasing resolution, indicating a mesoscale shear instability is present that acts to weaken the mean shear, and that it becomes much more active at higher resolution. The vertical flux of meridional momentum by mesoscale motions (difference between open symbols and closed symbols in Fig. 15d) also increases strongly in magnitude with resolution, but it is negative (upwards), which is an up-gradient flux, because the Gulf Stream is a region of positive shear in the meridional wind.

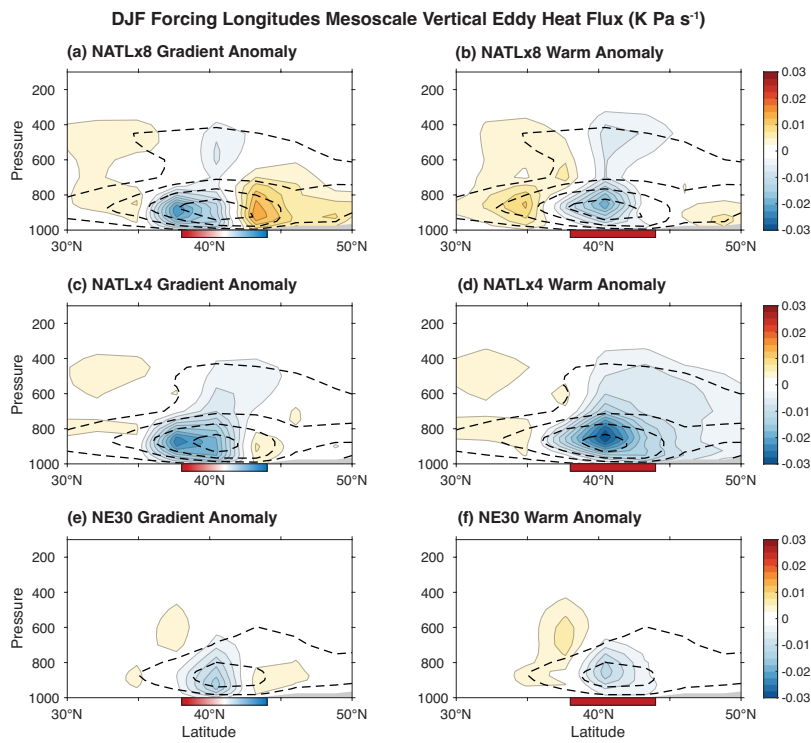
Returning to our discussion of the thermodynamic equation, the massive increases in vertical velocities with resolution has only a minor influence on the vertical EHF, because the increase in vertical velocities is primarily occurring at scales much smaller than the  $O(1000\text{ km})$  scale of most temperature anomalies. This can be seen by the similarity of the climatologies of the large-scale vertical EHF as resolution is changed (black contours in Fig. 16). There is a large increase in the mesoscale upward vertical EHF with resolution (black contours in Fig. 17); however, the mesoscale vertical EHF is an order of magnitude smaller than the large-scale vertical EHF. Here, as in Fig. 15, we are separating large-scale and mesoscale fluxes by switching the order of operations of computing the variance from the 6-hourly data and conservatively remapping to the  $\sim 200$ -km f19 grid, then using Reynold's decomposition.

While the contribution of mesoscale motions to the vertical EHF is small, it offers a potential explanation for what is driving the deep temperature anomaly in response to Gulf Stream SST anomalies, because the response of mesoscale vertical EHF to SST anomalies shows an upward heat flux extending into the upper troposphere in NATLx4 and NATLx8 (Fig. 17). While small in magnitude, this vertical EHF creates a direct link between the surface and the upper troposphere over the Gulf Stream. The large-scale vertical EHF response of opposite sign (Fig. 16) can be thought of as a response to the deep temperature anomaly and acts opposite to the mesoscale vertical EHF. However, both NATLx4 and NATLx8 show upward heat flux anomalies of comparable magnitude and vertical extent (Fig. 17), so why don't the NATLx4 simulations also show a deep

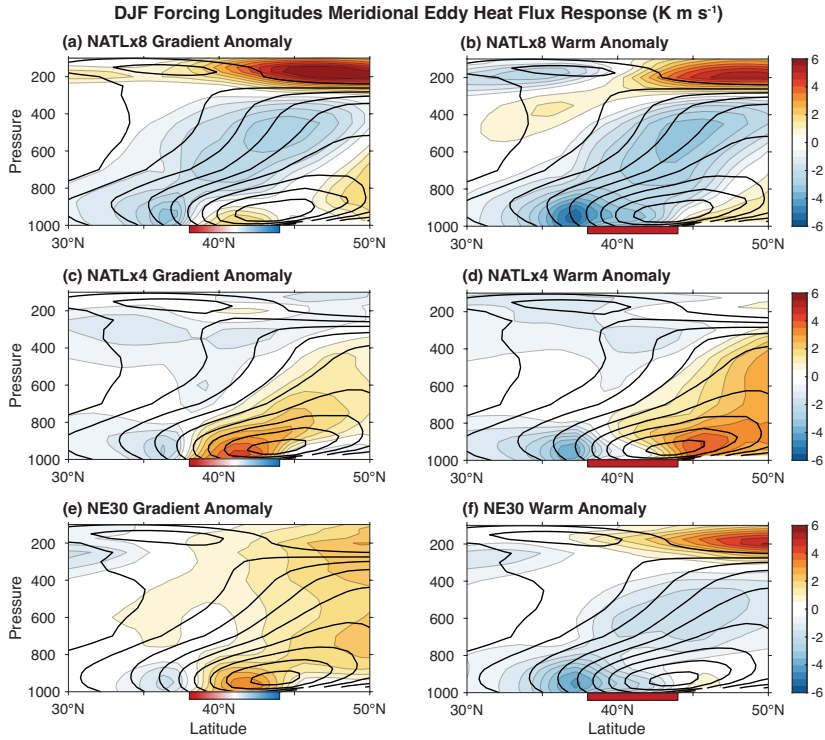




**Figure 16.** Same as Fig. 5, but for the DJF vertical eddy heat flux by large-scale motions, defined by the covariance of pressure velocity and temperature on scales greater than 200 km, computed as described in the text. Upward heat fluxes are negative. The contour interval for the climatology is  $0.2 \text{ K Pa s}^{-1}$ .



**Figure 17.** Same as Fig. 5, but for the DJF vertical eddy heat flux by mesoscale motions, defined by the covariance of pressure velocity and temperature on scales less than 200 km, computed as described in the text. Upward heat fluxes are negative. The contour interval for the climatology is  $0.03 \text{ K Pa s}^{-1}$ .



**Figure 18.** Same as Fig. 5, but for the DJF meridional eddy heat flux (total of large scale and mesoscale, the latter of which is negligible). The contour interval for the climatology is  $5 \text{ K m s}^{-1}$ .

663 temperature anomaly? A potential reason is that NATLx4 does not sufficiently distin-  
 664 guish between mesoscale and synoptic scale motions, so the upward heat fluxes from the  
 665 surface become part of the cyclone warm conveyor belts, which don't just move heat up-  
 666 ward but also poleward. This hypothesis is supported by the poleward and upward EHF  
 667 by large-scale (synoptic) motions in response to SST anomalies in NATLx4 (positive anom-  
 668 alies north of  $40^\circ\text{N}$  in Figs. 18c,d and negative anomalies north of  $40^\circ\text{N}$  in 16c,d), unlike  
 669 the EHF anomalies in NATLx8 and NE30-WARM (Figs. 18a,b,f and 16a,b,f). The dif-  
 670 ferences in meridional and vertical eddy heat fluxes are also reflected in North Atlantic  
 671 eddy kinetic energy (EKE), with EKE generally decreasing in NATLx8 versus shifting  
 672 poleward and increasing slightly in NATLx4 (Supporting Information Fig. S6).

673 Drawing on the analysis presented so far, we propose a potential explanation for  
 674 the difference in response between the NATLx8 and NATLx4 simulations: While effec-  
 675 tive buoyancy arguments (Jeevanjee & Romps, 2016) lead to an increase in magnitude  
 676 of resolved updrafts in both NATLx4 and NATLx8 relative to NE30, this ascent is more  
 677 concentrated within cold fronts (i.e., south-southeast of the cyclone center) in NATLx8  
 678 versus warm fronts (i.e., east-northeast of the cyclone center) in NATLx4 (Fig. 14). The  
 679 steep isentropic slopes of cold fronts lead to an efficient pathway for surface anomalies  
 680 to be communicated to the free troposphere by adiabatic motions, and the occurrence  
 681 of cold fronts within the sector of the cyclone with smaller meridional winds (relative to  
 682 warm fronts) means that there isn't a simultaneous poleward transport of these anom-  
 683 alies. This leads to a deep temperature response in NATLx8, whereas northward heat flux  
 684 within the warm sector of cyclones prevents this local warm anomaly from developing  
 685 in NATLx4. On the other hand, NE30-WARM also gets a deep temperature response,

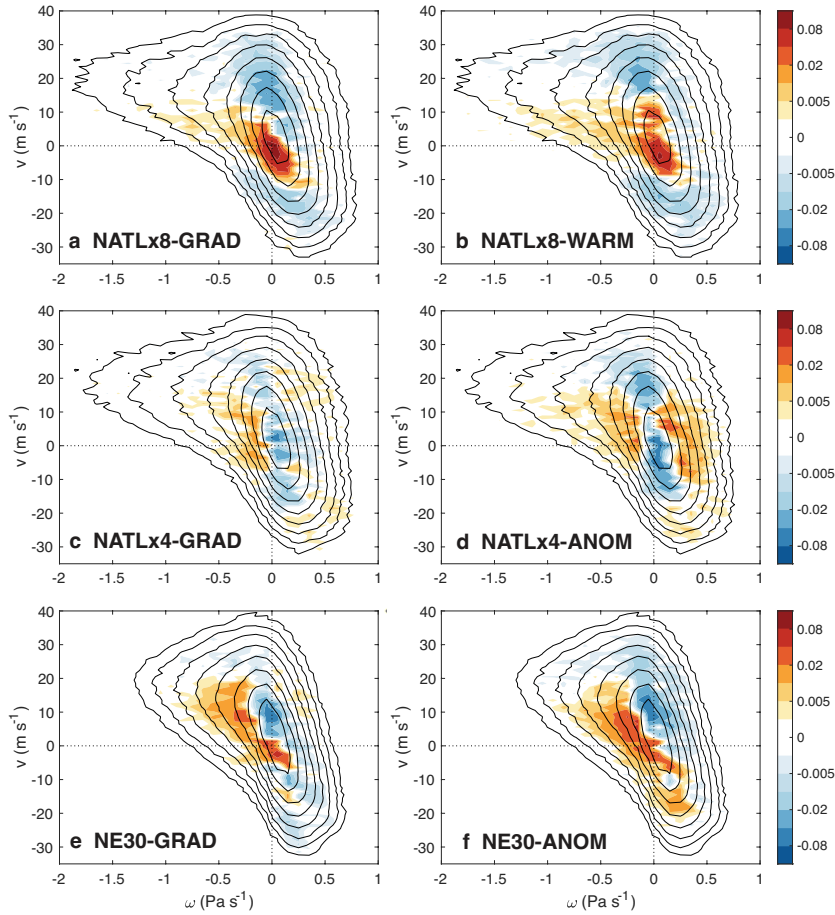
albeit weaker, which we speculate comes about via parameterized convection as opposed to the resolved ascent processes that govern the NATLx4 and NATLx8 responses.

This picture can be quantitatively supported by looking at changes in the covariance of vertical velocities and meridional winds ( $\omega'v'$ ; Fig. 15d), specifically at its response to SST anomalies. NATLx8 shows a decrease in the magnitude of  $\omega'v'$  in response to SST anomalies (i.e., the black triangle and black square are less negative than the black circle). This results from a contraction of the probability distribution for the meridional wind such that more ascent occurs with weakly positive meridional winds (e.g., cold front convection) and more descent occurs with weakly negative meridional winds (e.g., in the cold sector) (Fig. 19a,b). NATLx4 shows the opposite: an increase in the magnitude of  $\omega'v'$  in response to SST anomalies (Fig. 15c). While it shows a similar shift of strong ascent towards conditions with weaker meridional winds (i.e., from the warm front to the cold front) (Fig. 19c,d), it shows completely different anomalies in the weak ascent and descent parts of the joint probability distribution of  $\omega$  and  $v$ , such that overall it shows a strengthening of the existing covariance between vertical and meridional winds more than it shows a shift in the meridional winds at which ascent and descent are occurring. Notably, NATLx4-WARM in particular shows a shift of descent from weakly negative  $v$  to weakly positive  $v$  (Fig. 19d), which is consistent with the ascending air becoming entrained in the poleward traveling warm conveyor belt (Browning et al., 1973), where it later descends. The response of  $\omega'v'$  in NE30 is positive like in NATLx8 (Fig. 15c), but the response of the joint probability distribution of vertical and meridional wind looks different again, with a shift towards more upward and equatorward winds throughout the distribution (Fig. 19e,f).

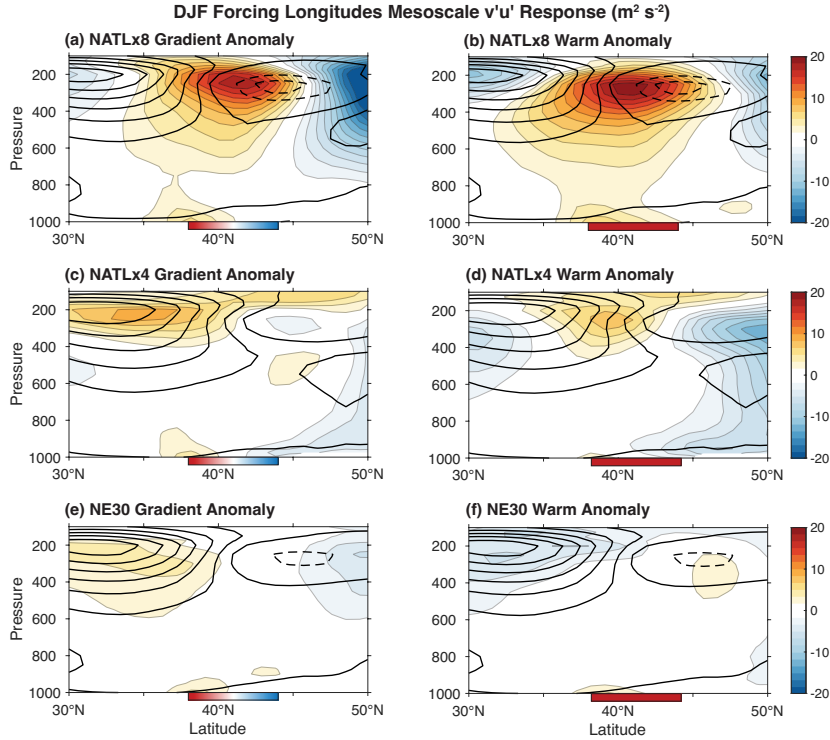
It is not just the transient-eddy heat flux responses that show large differences with resolution. Many transient-eddy fluxes show large differences in the response to idealized SST anomalies with resolution. A notable example is the meridional flux of zonal momentum by transient eddies (Fig. 20). NATLx8 shows strong poleward anomalies in the eddy momentum flux in response to both SST anomalies, which would help to explain the strong poleward shift of the jet in these simulations (Fig. 6). It is also consistent with the negative anomalies in poleward EHF (Fig. 18) and the strong positive anomalies to the north (not shown), which from an Eliassen-Palm flux perspective should be associated with a convergence of zonal momentum, as is seen at  $\sim 45^\circ\text{N}$ . The lower resolution simulations show much weaker anomalies in the meridional flux of zonal momentum by transient eddies. However, as with the thermodynamic analysis, it is difficult to disentangle the causality, i.e., whether the eddy fluxes of zonal momentum are an important reason for the large-scale circulation response or are themselves a result of the large-scale circulation response is challenging to parse out. One hint that this could be a causal factor in the response is that there is a large upward anomaly in mesoscale  $\omega'u'$  (Fig. 15c) at the south edge of the  $v'u'$  anomaly. Future work should investigate the momentum budget and in particular the strength of the eddy momentum flux feedback in this model configuration, because this feedback has been suggested to get stronger with higher resolution (Hardiman et al., 2022), and Fig. 20 provides some preliminary evidence of this.

## 4 Conclusions and Discussion

Our results show a large ( $\sim 2$  hPa  $(^\circ\text{C})^{-1}$ ) East-Atlantic-intensified positive NAO-like response to warm SST anomalies south of the Gulf Stream SST front in a variable-resolution version of CAM6 with 14-km regional grid refinement over the North Atlantic. This response is weaker and has a different spatial structure in lower resolution simulations, including in simulations with 28-km regional grid refinement over the North Atlantic, corresponding to the resolution used in many previous high-resolution modeling efforts (Haarsma et al., 2016; Chang et al., 2020). The key circulation responses to SST anomalies are tabulated across the different resolutions of our simulations in Table 1. The



**Figure 19.** Normalized bivariate probability distributions of 6-hourly instantaneous meridional wind  $v$  and vertical pressure velocity  $\omega$  within the Gulf Stream forcing region during DJF, on the model level with average pressure of 610 hPa. Contours show the climatology (REF), with contour intervals [0.005 0.01 0.02 0.04 0.08 0.16 0.32 0.64]. Shading shows the response to (left) SST gradient anomalies in the Gulf Stream (GRAD–REF) and (right) warm SST anomalies in the Gulf Stream (WARM–REF) on a log scale. 3 different configurations of CAM-SE are shown: (a),(b) NATLx8, with 14-km resolution in the North Atlantic, (c),(d) NATLx4, with 28-km resolution in the North Atlantic, and (e),(f) NE30, with global 111-km resolution. This analysis is based on data that has been regridded to the 100-km f09 analysis grid, such that it does not capture the magnitude of the strongest updrafts found in NATLx4 and NATLx8.



**Figure 20.** Same as Fig. 5, but for the DJF meridional eddy flux of zonal momentum (total of large scale and mesoscale, the latter of which is negligible). The contour interval for the climatology is  $8 \text{ m}^2 \text{ s}^{-2}$ .

738 larger circulation response in NATLx8 is particularly apparent in the negative SLP and  
 739 Z300 anomalies and in the zonal-mean jet shift (i.e., [U300]). The differences we find in  
 740 the large-scale circulation response result entirely from differences in horizontal atmo-  
 741 spheric resolution, because the same  $1^\circ$  resolution SSTs are specified at each atmospheric  
 742 resolution, and the vertical resolution is kept constant. Investigating how these responses  
 743 are modified by higher vertical resolution or ocean resolution is a promising avenue for  
 744 future research.

#### 745 4.1 Comparison with Observations

746 Given that our results are entirely based on a single atmospheric model (CAM6),  
 747 it is important to validate the response found in the high resolution simulations against  
 748 observations. We chose the Gulf Stream SST forcing region for our simulations based  
 749 on the observational analysis of S. M. Wills et al. (2016), making this study the clear-  
 750 est reference point. For a peak SST anomaly amplitude of  $1^\circ\text{C}$  in this region, they find  
 751 a 1000-hPa geopotential height response of  $\sim 14$  meters, corresponding to an SLP response  
 752 of  $\sim 1.7$  hPa at a near-surface density of  $1.25 \text{ kg m}^{-3}$ . This is in good agreement with  
 753 the  $1.9 \text{ hPa } (^\circ\text{C})^{-1}$  found in our NATLx8 simulations (Table 1), especially considering  
 754 that in the observational composite the SSTs only have a peak amplitude of  $1^\circ\text{C}$  and the  
 755 average over the Gulf Stream region is lower than this. However, the spatial pattern of  
 756 the response is quite different between NATLx8-WARM and the observational analogue  
 757 of S. M. Wills et al. (2016). Where NATLx8-WARM shows a weak high over the mid-  
 758 latitude North Atlantic and a strong low over the Norwegian Sea, the observational analogue  
 759 shows a weak low over the Gulf Stream, a strong high over the subpolar North At-

**Table 1.** Responses of key processes and circulation indices to GRAD and WARM SST anomalies for each resolution. The largest response for each SST anomaly is highlighted in bold.

Index	NATLx8 (14 km)		NATLx4 (28 km)		NE30 (110 km)	
	GRAD	WARM	GRAD	WARM	GRAD	WARM
EOF-1 NAO <sup>a</sup>	0.41	<b>0.44</b>	<b>0.52</b>	0.42	-0.12	0.35
EOF-2 <sup>a</sup>	<b>0.21</b>	<b>0.61</b>	0.15	-0.15	0.13	0.14
SLP (hPa) <sup>b</sup>	2.0, <b>-3.5</b>	1.1, <b>-3.9</b>	<b>2.4</b> , -2.5	<b>2.3</b> , -2.9	1.1, -0.72	1.5, -2.2
Z300 (m) <sup>b</sup>	30, <b>-41</b>	24, <b>-47</b>	<b>33</b> , -28	<b>28</b> , -29	14, -9.1	28, -24
U700 (m s <sup>-1</sup> ) <sup>b</sup>	<b>1.9</b>	<b>2.1</b>	<b>1.9</b>	1.8	0.49	1.4
[U300] (m s <sup>-1</sup> ) <sup>c</sup>	<b>2.0</b> , <b>-1.3</b>	<b>2.4</b> , <b>-1.8</b>	1.6, <b>-1.3</b>	0.89, -0.50	0.78, -0.63	1.2, -0.86
Near-Surface Convergence (10 <sup>-6</sup> s <sup>-1</sup> ) <sup>d</sup>	4.33	5.12	<b>5.04</b>	<b>5.35</b>	4.16	4.91
Precipitation (mm day <sup>-1</sup> ) <sup>d</sup>	<b>1.61</b>	1.89	<b>1.61</b>	<b>2.23</b>	1.28	1.87
Sfc. Heat Flux (W m <sup>-2</sup> ) <sup>d</sup>	88	91	<b>105</b>	<b>117</b>	93	113
$-\partial T/\partial y$ (°C (1000 km) <sup>-1</sup> ) <sup>d</sup>	0.76, <b>-0.92</b>	1.2, <b>-1.3</b>	<b>1.1</b> , -0.36	<b>1.4</b> , -0.71	0.56, -0.76	0.85, -1.0
$\theta_{300-800}$ <sup>e</sup>	<b>0.51</b>	<b>0.69</b>	0.21	0.15	0.20	0.36

<sup>a</sup> Mean anomalies in the PCs from Figure 7 in units of interannual standard deviations.<sup>b</sup> Max and min computed over 40°W-60°E in the Northern Hemisphere.<sup>c</sup> Max and min over 20-70°N of the values in Figure 6.<sup>d</sup> Max (and min for  $-\partial T/\partial y$ ) computed over 30-80°W and 30-50°N.<sup>e</sup> Max over 30-55°N of the 300-800 hPa vertical-average potential temperature anomaly.

760 lantic, and a weak low over Scandinavia and Northern Europe, more similar to the NATLx4-  
761 WARM response.

762 Rather than indicating a clear failure of the model, the differences in spatial pat-  
763 tern between the NATLx8-WARM response and the observational analogue (S. M. Wills  
764 et al., 2016) reflect differences in the associated SST pattern. The Gulf Stream SST in-  
765 dex analyzed by S. M. Wills et al. (2016) corresponds to variability in the latitude of the  
766 Gulf Stream (see also Famooss Paolini et al., 2022), with warm SSTs north of the Gulf  
767 Stream front corresponding to a more northerly Gulf Stream position. The SST pattern  
768 used in our simulations also includes warm SST anomalies south of the Gulf Stream front,  
769 which are found to be key to the large-scale circulation response (as indicated by the sim-  
770 ilarity of the responses in the GRAD and WARM experiments). Therefore, while the SST  
771 anomalies used in our simulations help to identify which aspects of the SST pattern mat-  
772 ter (i.e., the SSTs south of the Gulf Stream front in our simulations), they do not have  
773 a clear analogue in observed variability. For this reason, we plan to follow up on this work  
774 with simulations forced by SST patterns derived from observed variability, with the aim  
775 of making a clearer observational validation of the large-scale circulation response.

## 776 4.2 Mechanistic Understanding

777 The increased large-scale circulation response to Gulf Stream SST anomalies at high  
778 (14-km) resolution stems from an increase in resolved vertical motions within midlat-  
779 itude cyclones. The increase in vertical motion within midlatitude cyclones modifies transient-  
780 eddy fluxes of energy and momentum, especially their response to SST perturbations.  
781 In the highest (14-km) resolution simulations, mesoscale motions move anomalous heat  
782 from the surface into the free troposphere, where they help to sustain a temperature anomaly  
783 throughout the free troposphere over the Gulf Stream. Our results suggest that this is  
784 mostly due to convection in the cold sector, consistent with the mechanisms discussed  
785 by Vannière et al. (2017) in the context of an individual storm system.

786 Simulations with a lower resolution of 28 km, which is still high by climate mod-  
787 eling standards, show a qualitatively different response across many variables. Based on  
788 our analyses, we suggest that this is because at this resolution the upward heat trans-  
789 port by mesoscale circulations becomes part of the warm-conveyor belt, where warm moist  
790 air ascends and moves poleward in the warm sector of the cyclone. In this way the sig-  
791 nal from the surface anomalies doesn't ascend to the upper troposphere within the forc-  
792 ing region, but is instead moved poleward within the storm track. More work on the eddy-  
793 mean flow interactions in mesoscale-resolving models (especially those downstream in  
794 the East Atlantic) is needed to understand why this impact on the eddy heat flux does  
795 not translate into as large of an impact on the upper-tropospheric circulation. More gen-  
796 erally, a key next step in future work should be determining how the eddy-mean flow feed-  
797 back changes across this hierarchy of resolutions.

798 In determining how strong extratropical cyclones get and where they propagate and  
799 break (thus determining their influence on the mean flow), another mechanism that should  
800 be considered is that at high resolution the latent heating could become more concen-  
801 trated in fronts, which could cause more intense cyclones that propagate further pole-  
802 ward (Schemm, 2023). This mechanism, investigated by Schemm (2023) in idealized vari-  
803 able resolution simulations of the response to midlatitude SSTs, could operate in addi-  
804 tion to the mesoscale vertical heat transport mechanisms we investigate in this paper,  
805 and both could be important in understanding the large-scale circulation response. An-  
806 other important mechanism that we have not considered here is the influence of better  
807 orographic resolution on North Atlantic circulation regimes (Davini et al., 2022). We fo-  
808 cused on the eddy heat flux responses in the forcing region in this paper, because we found  
809 them to be important in explaining the differences between our 14-km and 28-km sim-  
810 ulations.



811 The differences in response between our 28-km and 14-km resolution simulations  
 812 suggests that increasing atmospheric resolution to resolve localized convective systems  
 813 embedded in cold fronts may lead to fundamental differences in how the atmosphere re-  
 814 sponds to midlatitude surface perturbations. Variable-resolution simulations, due to their  
 815 computational efficiency compared to mesoscale-resolving global simulations, offer a key  
 816 tool for understanding the upscale influence of mesoscale processes on large-scale dynam-  
 817 ics, a topic on which many open questions remain.

### 818 4.3 Implications

819 Our results have important implications for seasonal-to-decadal prediction, because  
 820 they suggest that models with higher resolution than is currently used have a larger at-  
 821 mospheric response to North Atlantic SST anomalies, which are predictable at lead times  
 822 of years to decades (Msadek et al., 2014; Meehl et al., 2014; S. G. Yeager et al., 2018;  
 823 Borchert et al., 2021; S. G. Yeager et al., 2023). If this response is indeed realistic and  
 824 can be reproduced with other SST patterns and within other models, then it suggests  
 825 that increasing the resolution of our seasonal-to-decadal prediction models to resolve frontal-  
 826 scale processes could lead to dramatic increases in skill in predicting decadal variations  
 827 in the atmospheric circulation and regional climate, e.g., for predicting precipitation in  
 828 Western Europe (Simpson et al., 2019). This is also relevant in the context of anthro-  
 829 pogenic climate change, where non-uniform warming features such as the North Atlantic  
 830 warming hole may elicit a larger forced atmospheric response, which may help to explain  
 831 model-observations discrepancies in long-term circulation trends (Blackport & Fyfe, 2022;  
 832 Vautard et al., 2023).

833 A larger response to North Atlantic SST anomalies also offers a potential resolu-  
 834 tion to the signal-to-noise paradox (Eade et al., 2014; Scaife & Smith, 2018; D. M. Smith  
 835 et al., 2020): current climate models are predicting something like the correct pattern  
 836 and phasing of atmospheric responses to SST anomalies but with too weak amplitude  
 837 (e.g., NE30-WARM response vs. NATLx8-WARM response in Fig. 3) such that the am-  
 838 plitude of the predictable signal is underestimated. Our results suggest that the signal-  
 839 to-noise paradox should get less severe as we increase the resolution of seasonal-to-decadal  
 840 prediction models to better resolve frontal processes and their role in communicating sur-  
 841 face anomalies into the upper troposphere. S. G. Yeager et al. (2023) have already found  
 842 evidence of this in other regions in a high resolution decadal prediction system using CESM  
 843 with a  $0.25^\circ$  atmospheric resolution and a  $0.1^\circ$  ocean resolution. While the  $2\times$  increase  
 844 in circulation responses from 110-km to 14-km resolution in our simulations (and less  
 845 than that for the NAO, see Table 1) is less than the factor or  $\sim 4$  by which some stud-  
 846 ies (e.g., D. M. Smith et al. (2020)) suggest the predictable NAO signal is underestimated,  
 847 there are several factors that could explain the remainder of the discrepancy: (1) there  
 848 is substantial uncertainty in the factor by which the predictable signal is underestimated,  
 849 due to the short observational record, and it could in fact be less than 4, (2) other SST  
 850 patterns besides the ones considered here could be more important for the difference in  
 851 the (e.g., NAO) response with resolution, and (3) the response could increase further with  
 852 even higher resolution or when coupled to a higher resolution ocean. It is also worth not-  
 853 ing that an increase in eddy-mean flow feedback at higher resolution, as suggested by  
 854 Hardiman et al. (2022), could increase the predictable signal independently of SST forc-  
 855 ing.

856 A larger atmospheric response to North Atlantic SST anomalies would mean a larger  
 857 feedback of the ocean state onto the further evolution of the SST anomalies. The details  
 858 of how this influences the atmosphere-ocean dynamics of decadal variability depends on  
 859 the sign and pattern of atmospheric response to realistic SST anomaly patterns, which  
 860 should be investigated in future work with mesoscale-resolving climate models. However,  
 861 if the atmospheric circulation response to realistic SST patterns is such as to further am-

862 plify those SST patterns, this could further increase the amplitude of predictable mid-  
 863 latitude signals.

## 864 5 Open Research

865 The CESM2.2 run scripts, grid files, and SST forcing files used to run our simu-  
 866 lations are available in a Zenodo repository (Jnglin Wills, 2024). The Zenodo repository  
 867 also contains model output used in the paper including (1) the DJF climatology of all  
 868 atmospheric fields for each simulation, (2) monthly-mean SLP for all months, (3) pentadal-  
 869 mean SLP and zonal wind at 850 hPa in the North Atlantic domain, and (4) climato-  
 870 logical covariances processed from 6-hourly model output needed for the separation of  
 871 fluxes into large-scale ( $> 200$  km) and mesoscale ( $< 200$  km) components as described  
 872 in the text. All output in the repository has been regridded to the f09 or f19 grids. Fi-  
 873 nally, the Zenodo repository also contains the MATLAB scripts needed to reproduce all  
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